

ASYMPTOTICS FOR ELASTIC CONTINUA - A



LSK (LYAPUNOV - SCHMIDT - KOITER) ASYMPTOTICS

• General asymptotic method to study motion of systems (discrete or continuous) near singular points. Here the method is applied to the equilibrium of conservative elastic systems.

• IDEA: Study the projection of equilibrium equations along the finite dimensional null space of the system's stability operator at critical point. This way the study of a large problem is reduced to the study of a nonlinear system of *m* equations, where *m* is the multiplicity of the stability operator's eigenvalue at the critical point.

• Method follows asymptotically equilibrium paths emerging from bifurcation points (simple or multiple) of perfect systems and determines their stability.

• Method also investigates the equilibrium and stability of imperfect systems, near critical points of their perfect counterparts, for small imperfection amplitudes.

• NOTE: Method is useful in determining post-bifurcation behavior and imperfection sensitivity in applications as well as in providing efficient numerical tools for finding solutions near the singular points of complex nonlinear systems.





PRELIMINARIES – FUNCTIONALS & THEIR DERIVATIVES

 $\mathcal{E}(u)$: functional (i.e. function of a function) $\mathcal{E} \in \mathbb{R}$ depends on $u \in U$

U: infinite dimensional (Hilbert) space (special case : finite (Euclidean) space $U=\mathbb{R}^n)$

 (u_1, u_2) : inner product of $u_i \in U$, $|| u || \equiv (u, u)^{1/2}$: norm of u

$$\begin{split} \mathcal{E}_{,u} &: \text{linear operator } U \longrightarrow \mathbb{R} \text{ is } Frechet \text{ derivative } \mathcal{E}_{,u} \left[\alpha u_1 + \beta u_2 \right] = \alpha \mathcal{E}_{,u} \left[u_1 \right] + \beta \mathcal{E}_{,u} \left[u_2 \right] \\ &\lim_{\|\delta u\| \to 0} | \ \mathcal{E}(u + \delta u) - \mathcal{E}(u) - \mathcal{E}_{,u} \left[\delta u \right] | \ / \| \ \delta u \| = 0 : \text{ Definition Frechet} \\ &\text{Special case : } \ \delta u = \epsilon v, \ \| \ v \| = 1, \ Gateau \text{ derivative } (Frechet \Longrightarrow Gateau) \\ &\lim_{\epsilon \to 0} | \ \mathcal{E}(u + \epsilon v) - \mathcal{E}(u) - \epsilon \mathcal{E}_{,u} \left[v \right] | \ / \epsilon = 0 : \text{ Definition Gateau} \\ &\mathcal{E}_{,u} \left[v \right] \equiv \mathcal{E}_{,u} \ v : \text{ used henceforth for notational simplicity} \end{split}$$





CALCULATING DERIVATIVES OF ANY ORDER

 $\mathcal{E}_{u} v = [\partial \mathcal{E}(u + \epsilon v) / \partial \epsilon]_{\epsilon=0}$: practical way to calculate first derivative

In analogy, the second derivative $\mathcal{E}_{,uu}$ is a bilinear operator $(\mathcal{E}_{,u}v)w$

 $(\mathcal{E}_{,uu} v)w = (\mathcal{E}_{,uu} w)v = [\partial^2 \mathcal{E}(u + \epsilon v + \zeta w)/\partial \epsilon \partial \zeta]_{\epsilon = \zeta = 0}$

The third derivative $\mathcal{E}_{,uuu}$ is a symmetric trilinear operator $((\mathcal{E}_{,uuu} v)w)z$

$$((\mathcal{E}_{,uuu}\,v)w)z = [\partial^3 \mathcal{E}(u + \epsilon v + \zeta w + \eta z)/\partial \epsilon \partial \zeta \partial \eta]_{\epsilon = \zeta = \eta = 0}$$

Using functional derivatives, Taylor series expansion of \mathcal{E} is found :

$$\mathcal{E}(u+\delta u) \approx \mathcal{E}(u) + \frac{1}{1!}\mathcal{E}_{,u}\,\delta u + \frac{1}{2!}(\mathcal{E}_{,uu}\,\delta u)\delta u + \frac{1}{3!}((\mathcal{E}_{,uuu}\,\delta u)\delta u)\delta u + \dots$$





CALCULATING FUNCTIONAL DERIVATIVES - EXAMPLES

$$\mathcal{E}(u) \equiv \int_{0} [\exp(w) + (2 + \cos(x))w^2 + 3(v_x^2 + v^2)]dx, \quad u = (v(x), w(x));$$

1

where:
$$w \in H^0[0,1]$$
, (i.e. $\int_0^1 w^2 < \infty$); $v \in H^1[0,1]$, (i.e. $\int_0^1 [v^2 + v,_x^2] dx < \infty$)

$$\begin{aligned} \mathcal{E}_{,u} \, u_1 &= \quad \frac{\partial}{\partial \epsilon_1} \left\{ \int_0^1 [\exp(w + \epsilon_1 w_1) + (2 + \cos x)(w + \epsilon_1 w_1)^2 + \\ &+ 3((v_{,x} + \epsilon_1 v_{1,x})^2 + (v + \epsilon_1 v_1)^2)] dx \right\}_{\epsilon_1 = 0} \\ \mathcal{E}_{,u} \, u_1 &= \quad \int_0^1 [\exp(w) w_1 + 2(2 + \cos x) w w_1 + 6(v_{,x} v_{1,x} + v v_1)] dx \end{aligned}$$

MEC563 – STABILITY OF SOLIDS: FROM STRUCTURES TO MATERIALS – LECTURE 3 Page 4





$$(\mathcal{E}_{,uu}\,u_1)u_2 = \frac{\partial^2}{\partial\epsilon_1\partial\epsilon_2} \left\{ \int_0^1 [\exp(w+\epsilon_1w_1+\epsilon_2w_2) + (2+\cos x)(w+\epsilon_1w_1+\epsilon_2w_2)^2 \right\}$$

$$+3((v_{,x}+\epsilon_{1}v_{1,x}+\epsilon_{2}v_{2,x})^{2}+(v+\epsilon_{1}v_{1}+\epsilon_{2}v_{2})^{2})]dx\bigg\}_{\epsilon_{1}=\epsilon_{2}=0}$$

$$(\mathcal{E}_{,uu}\,u_1)u_2 = \int_0^1 [\exp(w)w_1w_2 + 2(2+\cos x)w_1w_2 + 6(v_{1,x}v_{2,x}+v_1v_2)]dx$$

Similarly one obtains for higher derivatives...

$$((\mathcal{E}_{,uuu}\,u_1)u_2)u_3 = \int_{0}^{1} [\exp(w)w_1w_2w_3]dx$$

$$(((\mathcal{E}_{,uuuu}\,u_1)u_2)u_3)u_4 = \int_{0} [\exp(w)w_1w_2w_3w_4]dx$$









• About critical point u_c project solution increment Δu along null space \mathcal{N} and its complement \mathcal{N}^{\perp} .

- Solve equilibrium in \mathcal{N}^{\perp} and use $v(\xi, \Delta\lambda)$ to find equilibrium in \mathcal{N} from which you determine $\Delta\lambda$ as a function of ξ
- If $\Delta\lambda(\xi)$ is unique: limit point
- If $\Delta\lambda(\xi)$ is not unique: bifurcation





 $\mathcal{E}(u, \lambda)$: energy at displacement $u(\mathbf{x}) \in U$ and load $\lambda \ge 0$

 $\mathcal{E}(0,\lambda) = 0, \ \forall \lambda:$ zero energy at zero displacement

 $\mathcal{E}_{,u}(u,\lambda)\delta u = 0, \ \forall \ \delta u \in U:$ equilibrium statement

 $\mathcal{E}_{,u}(\overset{0}{u}(\lambda),\lambda)\delta u = 0, \ \forall \lambda; \quad \text{principal solution } \overset{0}{u}(\lambda), \ (\overset{0}{u}(0) = 0)$

 $\overset{0}{u}(\lambda)$ stable near $\lambda = 0$, i.e. min. eigenvalue of $\mathcal{E}_{,uu} (\overset{0}{u}(\lambda), \lambda) \equiv \mathcal{E}_{,uu}^{0}$ is $\overset{0}{\beta} > 0$

$$(\mathcal{E}_{,uu} \stackrel{0}{(u(\lambda),\lambda)} \delta u) \delta u \geq \stackrel{0}{\beta}(\lambda) \parallel \delta u \parallel^2, \ \stackrel{0}{\beta}(\lambda) > 0; \quad \exists \epsilon > 0, \ \forall \ \lambda \in [0,\epsilon]$$

NOTE: Unique & stable principal solution near zero load assumed (realistic structures)





 $\big[\mathcal{E}_{,uu}\,(\overset{0}{u}(\lambda),\lambda)(d\overset{0}{u}/d\lambda) + \mathcal{E}_{,u\lambda}\,(\overset{0}{u}(\lambda),\lambda)\big]\delta u = 0 \implies d\overset{0}{u}/d\lambda \text{ exists if } (\mathcal{E},\overset{0}{,uu})^{-1} \text{ exists}$

 $\mathcal{E}_{,uu}^{\ 0} \equiv \mathcal{E}_{,uu} \left(\overset{0}{u}(\lambda), \lambda \right) \text{ is positive definite, invertible, i.e. } \overset{0}{\beta}(\lambda) > 0, \text{ for } \lambda \in [0, \lambda_c)$

As load increases away from 0, the lowest load that $\overset{0}{u}$ cannot be continued is : λ_{c}

 $(\mathcal{E}_{,uu} (\overset{0}{u}(\lambda_c), \lambda_c)\overset{1}{u})\delta u = 0; \ \lambda_c: \text{ critical load, } \overset{1}{u}: \text{mode } (\parallel \overset{1}{u} \parallel = 1, \text{ assumed unique})$

In all directions orthogonal to null space \mathcal{N} of $\mathcal{E}^{c}_{,uu}$ operator still positive definite :

$$(\mathcal{E}_{,uu}^{c} \delta v) \delta v \geq \gamma \parallel \delta v \parallel^{2}, \quad \exists \gamma > 0, \quad \forall \delta v \in \mathcal{N}^{\perp}; \quad \text{where} : \mathcal{E}_{,uu}^{c} \equiv \mathcal{E}_{,uu} \left(\overset{0}{u} (\lambda_{c}), \lambda_{c} \right)$$
$$\mathcal{N} \equiv \{ u \in U \mid u = \mu \overset{1}{u}, \; \forall \mu \in \mathbb{R} \}, \; \mathcal{N}^{\perp} \equiv \{ v \in U \mid (v, \overset{1}{u}) = 0 \}; \quad (\mathcal{N} \oplus \mathcal{N}^{\perp} = U)$$





 $u = u_c + \xi \overset{1}{u} + v, \quad v \in \mathcal{N}^{\perp}, \quad \xi \in \mathbb{R}$

$$\mathcal{E}_{,v} \,\delta v = 0 \implies \mathcal{E}_{,u} \,(u_c + \xi u^1 + v, \lambda_c + \Delta \lambda) \delta v = 0; \quad \text{equilibrium in } \mathcal{N}^{\perp}$$

Expand about (u_c, λ_c) to find $v(\xi, \Delta \lambda)$, where :

$$v(\xi, \Delta \lambda) = \xi v_{\xi} + \Delta \lambda v_{\lambda} + \frac{1}{2!} \left[\xi^2 v_{\xi\xi} + 2\xi \Delta \lambda v_{\xi\lambda} + (\Delta \lambda)^2 v_{\lambda\lambda} \right] + \dots$$

 $\mathcal{E}_{\xi} = 0 \implies \mathcal{E}_{u} (u_c + \xi u^1 + v, \lambda_c + \Delta \lambda) u^1 = 0;$ equilibrium in \mathcal{N}

Expand about (u_c, λ_c) , using $v(\xi, \Delta \lambda)$, to find $\Delta \lambda(\xi)$

 $\Delta\lambda(\xi)$ is unique \implies limit load, $\Delta\lambda(\xi)$ is not unique \implies bifurcation





 $O(\xi)$: $(\mathcal{E}_{,uu}^{c} v_{\xi}) \delta v = 0 \implies v_{\xi} = 0, \quad (\mathcal{E}_{,uu}^{c} \text{ has unique eigenvalue } \overset{1}{u})$

- $O(\Delta\lambda) : (\mathcal{E}_{,uu}^{c} v_{\lambda} + \mathcal{E}_{,u\lambda}^{c}) \delta v = 0 \implies 0 \neq v_{\lambda} \in \mathcal{N}^{\perp}, \quad (\mathcal{E}_{,uu}^{c} \text{ invertible in } \mathcal{N}^{\perp})$
 - $O(\xi^2) : (\mathcal{E}_{,uu}^c v_{\xi\xi} + (\mathcal{E}_{,uuu}^c \overset{1}{u}) \overset{1}{u}) \delta v = 0 \implies 0 \neq v_{\xi\xi} \in \mathcal{N}^\perp$
- $O(\xi \Delta \lambda) : \quad (\mathcal{E}_{,uu}^{\ c} v_{\xi\lambda} + (\mathcal{E}_{,uuu}^{\ c} v_{\lambda} + \mathcal{E}_{,uu\lambda}^{\ c})^{1}_{u})\delta v = 0 \implies 0 \neq v_{\xi\lambda} \in \mathcal{N}^{\perp}$

 $O((\Delta\lambda)^2) : \quad (\mathcal{E}_{,uu}^{\ c} v_{\lambda\lambda} + (\mathcal{E}_{,uuu}^{\ c} v_{\lambda})v_{\lambda} + 2\mathcal{E}_{,uu\lambda}^{\ c} v_{\lambda} + \mathcal{E}_{,u\lambda\lambda}^{\ c})\delta v = 0 \implies 0 \neq v_{\lambda\lambda} \in \mathcal{N}^{\perp}$

Using above results in equilibrium along ${\mathcal N}$

$$0 = \Delta\lambda(\mathcal{E},_{u\lambda}^{c} \overset{1}{u}) + \frac{1}{2} [\xi^{2}((\mathcal{E},_{uuu}^{c} \overset{1}{u})\overset{1}{u})\overset{1}{u} + 2\xi\Delta\lambda((\mathcal{E},_{uuu}^{c} v_{\lambda} + \mathcal{E},_{uu\lambda}^{c})\overset{1}{u})\overset{1}{u} + (\Delta\lambda)^{2}((\mathcal{E},_{uuu}^{c} v_{\lambda})v_{\lambda} + 2\mathcal{E},_{uu\lambda}^{c} v_{\lambda} + \mathcal{E},_{u\lambda\lambda}^{c})\overset{1}{u}] + \dots$$











 $u = \overset{0}{u}(\lambda) + \xi \overset{1}{u} + v, \quad v \in \mathcal{N}^{\perp}, \quad \xi \in \mathbb{R}, \quad \xi \equiv (u - \overset{0}{u}, \overset{1}{u}):$ bifurcation amplitude

$$\mathcal{E}_{,v}\,\delta v = 0 \implies \mathcal{E}_{,u}\,(\overset{0}{u}(\lambda_c + \Delta\lambda) + \xi\overset{1}{u} + v(\xi, \Delta\lambda), \lambda_c + \Delta\lambda)\delta v = 0; \quad \text{equilibrium in } \mathcal{N}^{\perp}$$

Expand about (u_c, λ_c) to find $v(\xi, \Delta \lambda)$, where :

$$v(\xi, \Delta \lambda) = \xi v_{\xi} + \Delta \lambda v_{\lambda} + \frac{1}{2!} \left[\xi^2 v_{\xi\xi} + 2\xi \Delta \lambda v_{\xi\lambda} + (\Delta \lambda)^2 v_{\lambda\lambda} \right] + \dots$$
$$\mathcal{E}_{\xi} = 0 \implies \mathcal{E}_{u} \left(\overset{0}{u} (\lambda_c + \Delta \lambda) + \xi \overset{1}{u} + v(\xi, \Delta \lambda), \lambda_c + \Delta \lambda \right) \overset{1}{u} = 0; \quad \text{equilibrium in } \mathcal{N}$$

Expand about (u_c, λ_c) , using $v(\xi, \Delta \lambda)$, to find $\Delta \lambda(\xi)$ where :

$$\Delta\lambda(\xi) = \lambda_1\xi + \lambda_2\frac{\xi^2}{2!} + \lambda_3\frac{\xi^3}{3!} + \dots$$





$$O(\xi)$$
 : $(\mathcal{E}_{,uu}^{c} v_{\xi}) \delta v = 0 \implies v_{\xi} = 0, \quad (\mathcal{E}_{,uu}^{c} \text{ unique eigenvalue : } \overset{1}{u})$

$$O(\Delta\lambda) : \quad (\mathcal{E}_{,uu}^{c} v_{\lambda} + \frac{\mathcal{E}_{,uu}^{c} (du^{0}/d\lambda)_{c} + \mathcal{E}_{,u\lambda}^{c}}{\delta v}) \delta v = (\mathcal{E}_{,uu}^{c} v_{\lambda}) \delta v = 0 \implies v_{\lambda} = 0, \quad (\text{same})$$

$$O(\xi^2) : (\mathcal{E}_{uu}^{c} v_{\xi\xi} + (\mathcal{E}_{uuu}^{c} \overset{1}{u})^{1}_{u})\delta v = 0$$

$$O(\xi \Delta \lambda) : (\mathcal{E}_{,uu}^{c} v_{\xi\lambda} + (\mathcal{E}_{,uuu}^{c} (du^{0}/d\lambda)_{c} + \mathcal{E}_{,uu\lambda}^{c})^{1} \delta v = 0$$

$$O((\Delta\lambda)^2) : \left(\mathcal{E}_{,uu}^c v_{\lambda\lambda} + \left(\mathcal{E}_{,uuu}^c (du^0/d\lambda)_c\right)(du^0/d\lambda)_c + 2\mathcal{E}_{,uu\lambda}^c (du^0/d\lambda)_c + \mathcal{E}_{,uu\lambda}^c (du^0/d\lambda)_c + \mathcal{E}_{,uu\lambda}^c (du^0/d\lambda^2)_c\right) \delta v = (\mathcal{E}_{,uu}^c v_{\lambda\lambda}) \delta v = 0 \implies v_{\lambda\lambda} = 0, \quad \text{(same)}$$

NOTE A:
$$v(0, \Delta \lambda) = u - \overset{0}{u}(\lambda) = 0 \implies v_{\lambda} = v_{\lambda\lambda} = v_{\lambda\lambda\lambda} = \dots = 0$$

NOTE B: Highlighted terms λ derivatives of principal equilibrium : $\mathcal{E}_{,u}(\overset{0}{u}(\lambda),\lambda)\delta u = 0$









STABILITY – SIMPLE BIFURCATION CASE

 $(\mathcal{E}_{,uu} (\overset{0}{u}(\lambda), \lambda)\overset{0}{x}(\lambda))\delta u = \overset{0}{\beta}(\lambda)(\overset{0}{x}(\lambda), \delta u), \quad \text{stability of principal path}$

 ${}^{0}_{\beta}(\lambda)$: eigenvalue, ${}^{0}_{x}(\lambda)$: eigenvector of \mathcal{E}_{uu}^{0} ; $({}^{0}_{x}(\lambda), {}^{0}_{x}(\lambda)) = 1$

Evaluate at λ_c : $(\mathcal{E}_{uu}^{c} \overset{0}{x}(\lambda_c))\delta u = 0; \quad \overset{0}{\beta}(\lambda_c) = 0, \ (\overset{0}{x}(\lambda_c), \overset{0}{x}(\lambda_c)) = 1 \implies \overset{0}{x}(\lambda_c) = \overset{1}{u}$

Differentiate at λ_c : $((\mathcal{E}_{uuu}^c, (d\hat{u}/d\lambda)_c + \mathcal{E}_{uu\lambda}^c)^1 + \mathcal{E}_{uu\lambda}^c, (d\hat{u}/d\lambda)_c)\delta u = (d\beta/d\lambda)_c(\hat{u}, \delta u)$

Substitute : $\delta u = \overset{1}{u}$, recall : $(\overset{1}{u}, \overset{1}{u}) = 1 \implies ((d\mathcal{E}_{,uu} / d\lambda)_c \overset{1}{u}) \overset{1}{u} = (\overset{0}{d\beta} / d\lambda)_c$

Assumption : $(d\beta/d\lambda)_c < 0$; (recall : $\beta(\lambda) > 0$, $\forall \lambda \in [0, \lambda_c)$, holds in most applications)

Eigenvalues and eigenvectors of stability operator along principal path





STABILITY – SIMPLE BIFURCATION CASE

 $(\mathcal{E}_{,uu} \left(\stackrel{0}{u} (\lambda_c + \Delta \lambda(\xi) + \xi \stackrel{1}{u} + v(\xi, \Delta \lambda(\xi)), \lambda_c + \Delta \lambda(\xi)) x(\xi) \right) \delta u = \beta(\xi) (x(\xi), \delta u), \parallel x(\xi) \parallel = 1$

Use : $\beta(\xi) = \xi \beta_1 + \frac{\xi^2}{2} \beta_2 + \dots, \ x(\xi) = x_0 + \xi x_1 + \frac{\xi^2}{2} x_2 + \dots, \text{ expand about : } (u_c, \lambda_c)$

$$O(1)$$
 : $(\mathcal{E}_{uu}^{c} x_{0})\delta u = 0, \ (x_{0}, x_{0}) = 1; \implies x_{0} = \overset{1}{u}$

Assume :
$$((\mathcal{E},_{uuu}^{c} \overset{1}{u})\overset{1}{u})\overset{1}{u} \neq 0$$
, substitute : $\delta u = \overset{1}{u}, \ \lambda_{1} = -\frac{1}{2} \frac{((\mathcal{E},_{uuu}^{c} \overset{1}{u})\overset{1}{u})\overset{1}{u})\overset{1}{u}}{((d\mathcal{E},_{uu}/d\lambda)_{c}\overset{1}{u})\overset{1}{u}}$

$$O(\xi) : ((\mathcal{E}_{,uuu}^{c} (\lambda_{1}(d\hat{u}/d\lambda)_{c} + \hat{u}))\hat{u} + \lambda_{1}\mathcal{E}_{,uu\lambda}^{c} \hat{u} + \mathcal{E}_{,uu}^{c} x_{1})\delta u = \beta_{1}(\hat{u},\delta u), \Longrightarrow$$
$$\beta_{1} = \lambda_{1}((d\mathcal{E}_{,uu}/d\lambda)_{c}\hat{u})\hat{u}^{1} + ((\mathcal{E}_{,uuu}^{c} \hat{u})\hat{u})\hat{u} = \lambda_{1}[-((d\mathcal{E}_{,uu}/d\lambda)_{c}\hat{u})\hat{u}]$$





STABILITY – SIMPLE BIFURCATION CASE

Assume : $((\mathcal{E},_{uuu}^{c} \overset{1}{u})\overset{1}{u})\overset{1}{u} = 0$, substitute : $\delta u = \delta v$, $\lambda_1 = \beta_1 = 0$

 $O(\xi) : ((\mathcal{E}_{uuu}^{c} \overset{1}{u})^{1}_{u} + \mathcal{E}_{uu}^{c} x_{1})\delta v = 0, \quad (x_{1}, \overset{1}{u}) = 0, \implies x_{1} = v_{\xi\xi}$

MEC563 – STABILITY OF SOLIDS: FROM STRUCTURES TO MATERIALS – LECTURE 3 Page 17





ENERGY – SIMPLE BIFURCATION CASE

$$\Delta \mathcal{E} \equiv \mathcal{E}(\overset{0}{u}(\lambda_{c} + \Delta \lambda) + \xi \overset{1}{u} + v(\xi, \Delta \lambda), \lambda_{c} + \Delta \lambda) - \mathcal{E}(\overset{0}{u}(\lambda_{c} + \Delta \lambda), \lambda_{c} + \Delta \lambda)$$

Using previous results from asymptotic expansion about : (u_c, λ_c)

$$\Delta \mathcal{E} = \begin{cases} \frac{\xi^3}{6} \lambda_1 [-(d\beta/d\lambda)_c] + O(\xi^4) & \text{for asymmetric bifurcation}: ((\mathcal{E},_{uuu}^c \overset{1}{u}) \overset{1}{u}) \overset{1}{u} \neq 0 \\ \frac{\xi^4}{8} \lambda_2 [-(d\beta/d\lambda)_c] + O(\xi^5) & \text{for symmetric bifurcation}: ((\mathcal{E},_{uuu}^c \overset{1}{u}) \overset{1}{u}) \overset{1}{u} = 0 \end{cases}$$

Comparing energy of principal & bifurcated paths for same load







MEC563 – STABILITY OF SOLIDS: FROM STRUCTURES TO MATERIALS – LECTURE 3 Page 19









• LSK asymptotic expansion for simple bifurcation point, reduces infinite dimensional problem to the study of one dimensional problem (projection of equilibrium on the critical operator's null space).

• Principal branch changes stability at critical load (from stable at lower loads to unstable).

• For transcritical (asymmetric) bifurcation, supercritical portion of path is stable, subcritical portion of path is unstable.

- For symmetric bifurcations, supercritical paths are stable, subcritical paths are unstable.
- Stable paths have, for a given load, less energy than their neighboring unstable paths

NOTE: General case asymptotics obtained here are similar to results of the rigid T model!





ELASTICA EXAMPLE – MODEL SETTING & ENERGY



$$\theta(s) \in H_0^1[0, L], \quad \text{(i.e. } \int_0^L [(d\theta/ds)^2 + \theta^2]^{1/2} ds < \infty, \ \theta(0) = \theta(L) = 0)$$

inner product : $(\theta_1, \theta_2) \equiv (1/L) \int_0^L [\theta_1(s)\theta_2(s)] ds$





ELASTICA EXAMPLE – GOVERNING EQUATIONS

$$\mathcal{E}_{,u} \,\delta u = \int_{0}^{L} \left[EI\left(\frac{d\theta}{ds}\right) \left(\frac{d\delta\theta}{ds}\right) - \lambda \sin(\theta) \delta\theta \right] ds = 0, \implies \text{(integrate by parts)}$$
$$\mathcal{E}_{,u} \,\delta u = \int_{0}^{L} \left[-EI\left(\frac{d^{2}\theta}{ds^{2}}\right) - \lambda \sin(\theta) \right] \delta\theta ds + \left[EI\frac{d\theta}{ds} \delta\theta \right]_{s=L} - \left[EI\frac{d\theta}{ds} \delta\theta \right]_{s=0} = 0$$

 $\text{Recall}: \ \delta\theta(s) \text{ arbitraty function} \in H^1_0[0,L], \ \delta\theta(0) = \delta\theta(L) = 0, \quad \Longrightarrow$

 $\frac{d^2\theta}{ds^2} + \frac{\lambda}{EI}\sin\theta = 0, \text{ Elastica governing equation (Euler 1744);} \quad \theta(0) = \theta(L) = 0,$

 $\overset{0}{u}(\lambda) = \overset{0}{\theta}(s,\lambda) = 0 \quad \forall \lambda, \text{ principal solution}$





ELASTICA EXAMPLE – STABILITY OF PRINCIPAL PATH

$${}^{0}_{\beta}(\lambda) = \min_{\left(\|\delta u\|=1\right)} \left[\mathcal{E}_{,uu} \left({}^{0}_{u}(\lambda), \lambda \right) \delta u \right] = \int_{0}^{L} \left[EI \left(\frac{d(\delta \theta)}{ds} \right)^{2} - \lambda (\delta \theta)^{2} \right] ds, \ \forall \delta \theta \in H_{0}^{1}[0, L]$$

$$\delta\theta = \sum_{n=1}^{\infty} \left[\delta\theta_n \sin\left(\frac{n\pi s}{L}\right) \right], \ \left(\frac{1}{2} \sum_{n=1}^{\infty} \left(\delta\theta_n\right)^2 = \parallel \delta\theta \parallel = 1 \right), \text{ Fourier representation of } \delta\theta$$

$${}^{0}_{\beta}(\lambda) = \min_{(\delta\theta_n)} \frac{L}{2} \sum_{n=1}^{\infty} \left\{ (\delta\theta_n)^2 \left[EI\left(\frac{n\pi}{L}\right)^2 - \lambda \right] \right\} = L \left[EI\left(\frac{\pi}{L}\right)^2 - \lambda \right], \frac{d^{0}_{\beta}(\lambda)}{d\lambda} = -L < 0$$

Stability of principal branch changes at : $\lambda_c = EI\left(\frac{\pi}{L}\right)^2$, $\overset{0}{\beta}(\lambda) > 0 \ \forall \lambda \in [0, \lambda_c)$





ELASTICA EXAMPLE – CRITICAL LOAD AND MODE

$$(\mathcal{E}_{,uu} (\stackrel{0}{u} (\lambda_c), \lambda_c) \stackrel{1}{u}) \delta u = \int_0^L \left[EI \left(\frac{d\theta}{ds} \right) \left(\frac{d\delta\theta}{ds} \right) - \lambda_c \cos(\stackrel{0}{\theta}) \stackrel{1}{\theta} \delta \theta \right] ds = 0, \implies \text{(integrate by parts)}$$
$$(\mathcal{E}_{,uu} (\stackrel{0}{u} (\lambda_c), \lambda_c) \stackrel{1}{u}) \delta u = \int_0^L \left[-EI \left(\frac{d^2\theta}{ds^2} \right) - \lambda_c \stackrel{1}{\theta} \right] \delta \theta ds + \left[EI \frac{d\theta}{ds} \delta \theta \right]_{s=L} - \left[EI \frac{d\theta}{ds} \delta \theta \right]_{s=0} = 0$$

 $\text{Recall}: \ \delta\theta(s) \text{ arbitraty function} \in H^1_0[0,L]; \ \delta\theta(0) = \delta\theta(L) = 0, \quad \Longrightarrow \quad$

$$\frac{d^2\hat{\theta}}{ds^2} + \frac{\lambda_c}{EI}\hat{\theta} = 0, \quad \hat{\theta}(0) = \hat{\theta}(L) = 0; \quad \text{critical load}: \ \lambda_c, \quad \text{mode}: \ \hat{\theta} \ \left(\frac{1}{L}\int_0^L (\hat{\theta})^2 ds = 1\right)$$

 $\lambda_c = EI(\pi/L)^2, \quad \stackrel{1}{\theta}(s) = \sqrt{2}\sin(\pi s/L); \quad (\aleph_0 \text{ eigenvalues}: EI(n\pi/L)^2, \text{ modes}: \sqrt{2}\sin(n\pi s/L))$





ELASTICA EXAMPLE – LSK ASYMPTOTICS

$$\mathcal{E}_{u\lambda}^{c} \overset{1}{u} = -\int_{0}^{L} \sin(\overset{0}{\theta}(s))\overset{1}{\theta}(s)ds = 0, \quad ((\mathcal{E}_{uuu}^{c} \overset{1}{u})\overset{1}{u})\overset{1}{u} = \lambda_{c}\int_{0}^{L} \sin(\overset{0}{\theta}(s))[\overset{1}{\theta}(s)]^{3}ds = 0$$

Recall:
$$\lambda_2 = -\frac{1}{3} [(((\mathcal{E},_{uuuu}^c \overset{1}{u})\overset{1}{u})\overset{1}{u})\overset{1}{u})\overset{1}{u} + 3((\mathcal{E},_{uuu}^c v_{\xi\xi})\overset{1}{u})\overset{1}{u}]/((d\mathcal{E},_{uu}/d\lambda)_c \overset{1}{u})\overset{1}{u}]$$

$$(((\mathcal{E}_{,uuuu}^{c}\overset{1}{u})\overset{1}{u})\overset{1}{u})\overset{1}{u})\overset{1}{u} = \lambda_{c} \int_{0}^{L} \cos(\overset{0}{\theta}(s))[\overset{1}{\theta}(s)]^{4} ds = \frac{3}{2}L\lambda_{c} = \frac{3}{2}EI\frac{\pi^{2}}{L}$$

$$((\mathcal{E}_{,uuu}^{c}v_{\xi\xi})^{1}u)^{1}u = \lambda_{c}\int_{0}^{L}\sin(\overset{0}{\theta}(s))\theta_{\xi\xi}(s)[\overset{1}{\theta}(s)]^{2}ds = 0$$

$$\left((d\mathcal{E}_{,uu}/d\lambda)_{c}^{1}\hat{u}\right)^{1}\hat{u} = -\int_{0}^{L}\cos(\overset{0}{\theta}(s))[\overset{1}{\theta}(s)]^{2}ds = -L = \frac{d\overset{0}{\beta}(\lambda)}{d\lambda} \quad \text{(checks independently!)}$$

 $\lambda_2 = -\frac{1}{3} (\frac{3}{2}L\lambda_c)/(-L) = \frac{\lambda_c}{2} > 0, \quad \text{stable, supercritical, symmetric, bifurcation}$





ELASTICA EXAMPLE – REVIEW

• Euler's elastica has a trivial principal solution (straight configuration) that changes stability at the critical load.

- Critical load and mode depend on boundary conditions.
- Bifurcation at critical load is a simple (unique eigenmode), symmetric bifurcation.
- Principal path changes stability at critical load.
- Bifurcated (symmetric) path is supercritical near the critical load, which implies that it is stable.
- Asymptotic results are confirmed by exact solution (in terms of elliptic integrals).

Exact:
$$\begin{split} &L(\frac{\lambda}{EI})^{1/2} = 2K(\sin(\frac{\alpha}{2})) \\ &K(k) = \int_{0}^{1} \frac{dy}{[(1-y^2)(1-k^2y^2)]^{1/2}} \end{split} \text{Asymptotic:} \ \lambda = EI(\frac{\pi}{L})^2 [1 + \frac{1}{2}\sin^2(\frac{\alpha}{2}) + \dots] \end{split}$$