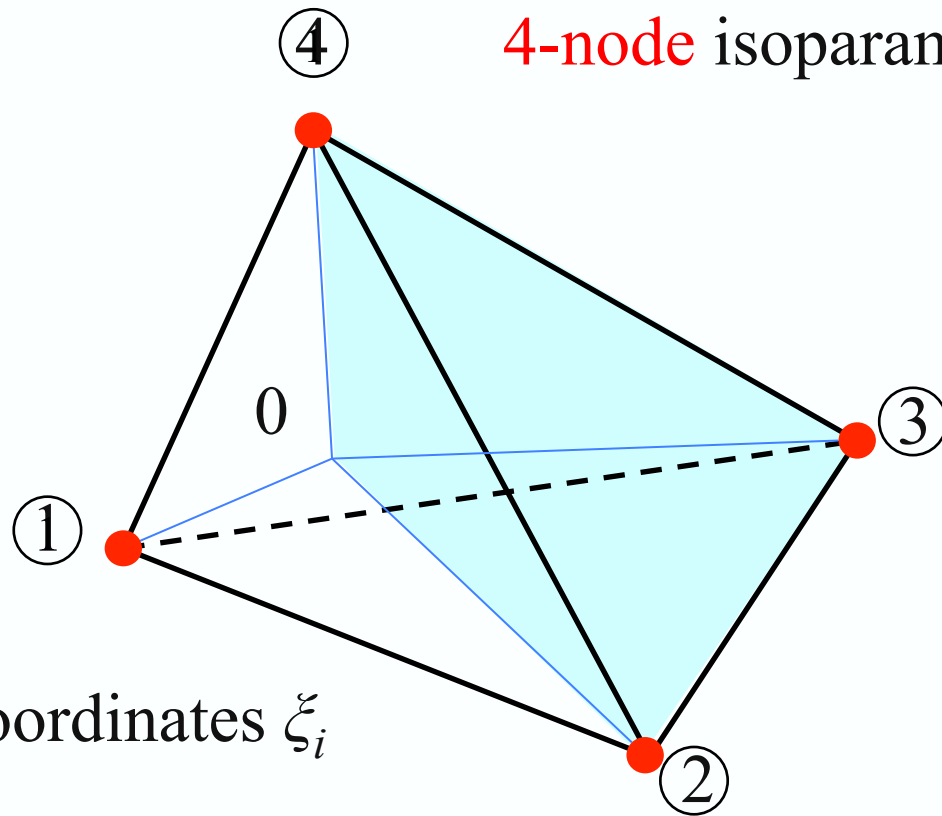


TOPICS COVERED IN THIS LECTURE

1. ISOPARAMETRIC ELEMENTS IN 3D (TETRAHEDRA, BRICKS ETC.)
2. TORSION OF A BAR IN 3D AND MOTIVATION FOR STRESS-BASED APPROACH
3. STRESS-BASED FEM; PRINCIPLE OF MINIMUM COMPLEMENTARY ENERGY
4. FEM DISCRETIZATION USING COMPLEMENTARY ENERGY

ISOPARAMETRIC ELEMENTS IN 3D

4-node isoparametric tetrahedron



$$u_j(\boldsymbol{\xi}) = \sum_{I=1}^4 N_I(\boldsymbol{\xi}) U_j^I$$

$$x_j(\boldsymbol{\xi}) = \sum_{I=1}^4 N_I(\boldsymbol{\xi}) x_j^I$$

Coordinates ξ_i

$$\xi_1 = V_{0234}/V_{1234} - V_{0234} : \text{Volume } 0234 \dots$$

$$\xi_2 = V_{0134}/V_{1234}$$

$$\xi_3 = V_{0124}/V_{1234}$$

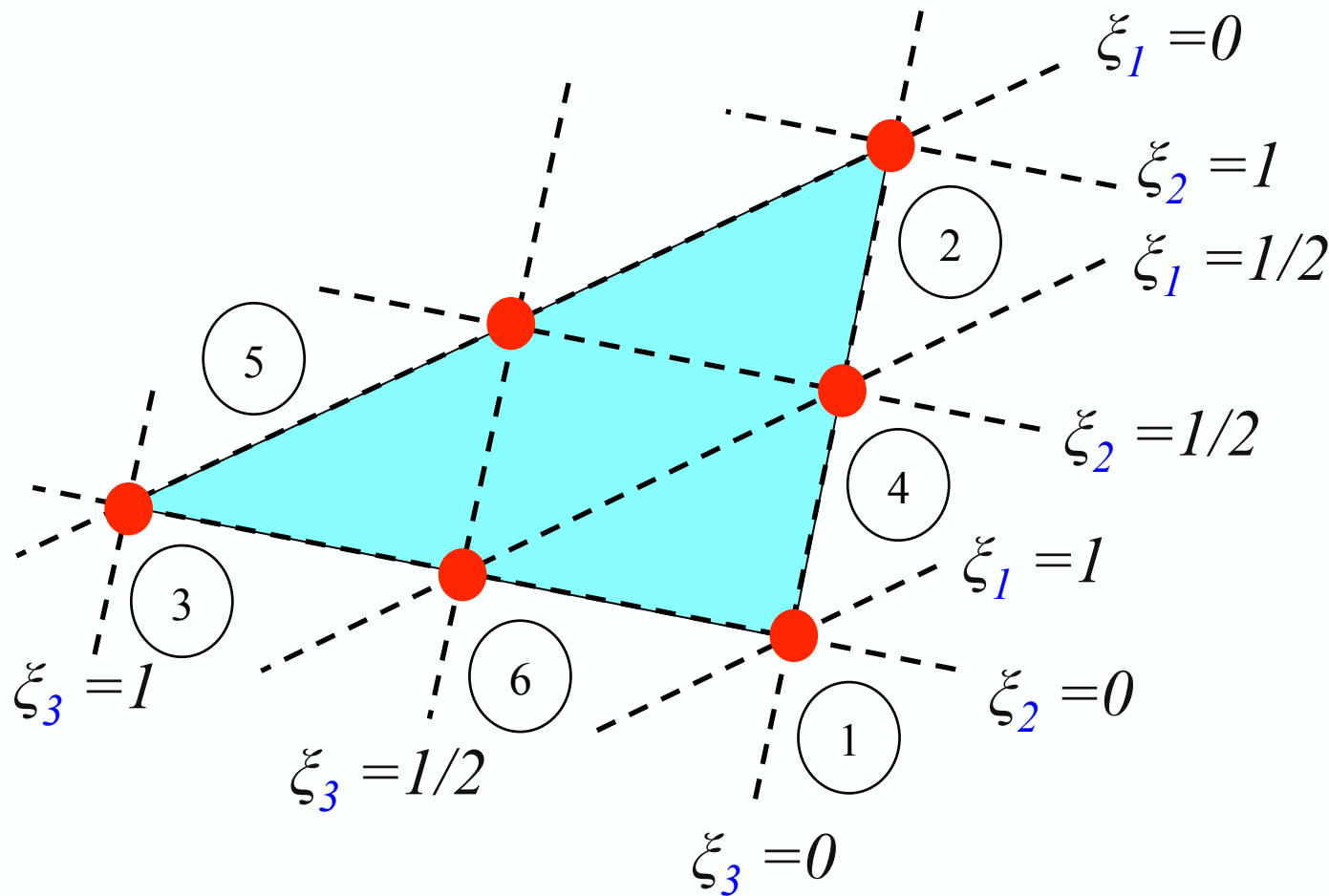
$$\xi_4 = V_{0123}/V_{1234}$$

$$N_I(\boldsymbol{\xi}^J) = \delta_{IJ}; \quad (I, J = 1, \dots, 4)$$

$$N_I(\boldsymbol{\xi}) = \xi_I; \quad (I = 1, \dots, 4)$$

Note: $\xi_1 + \xi_2 + \xi_3 + \xi_4 = 1$, patch test automatically **satisfied**

Shape functions of node I are products of equations avoiding that node



$$N_1 = \xi_1(2\xi_1 - 1)$$

$$N_2 = \xi_2(2\xi_2 - 1)$$

$$N_3 = \xi_3(2\xi_3 - 1)$$

$$N_4 = 4 \xi_1 \xi_2$$

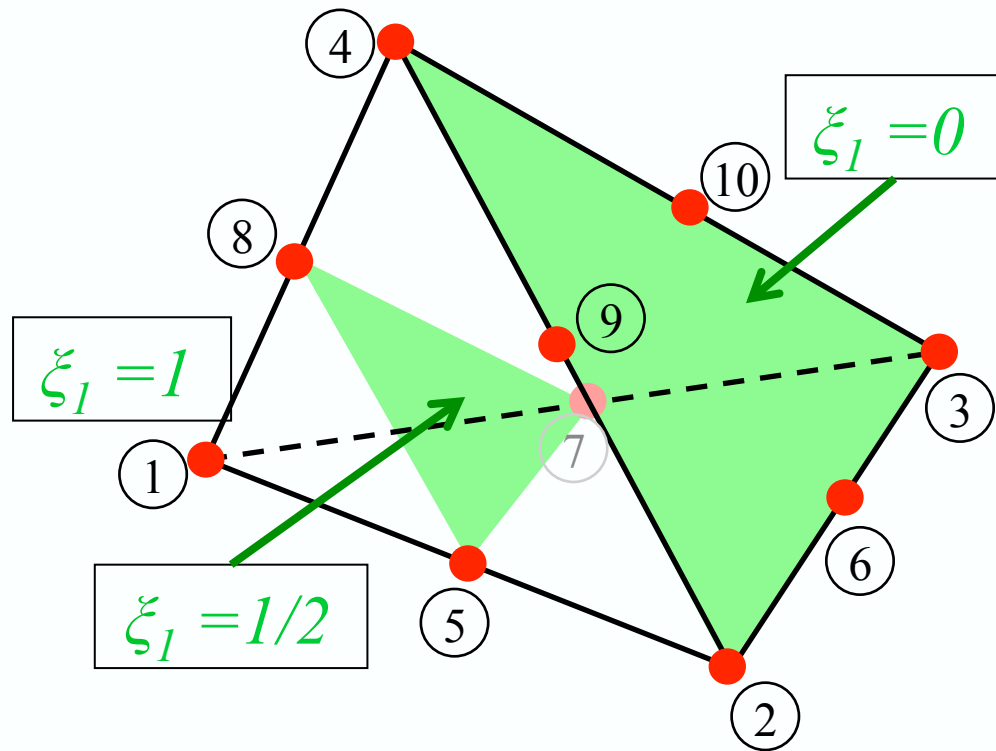
$$N_5 = 4 \xi_2 \xi_3$$

$$N_6 = 4 \xi_3 \xi_1$$

Triangular coordinates satisfy: $\xi_1 + \xi_2 + \xi_3 = 1$

10-node isoparametric tetrahedron in 3D similarly to 6-node triangle in 2D

Tetrahedral coordinates: $\xi_1 + \xi_2 + \xi_3 + \xi_4 = 1$



$$N_1 = \xi_1(2\xi_1 - 1)$$

$$N_2 = \xi_2(2\xi_2 - 1)$$

$$N_3 = \xi_3(2\xi_3 - 1)$$

$$N_4 = \xi_4(2\xi_4 - 1)$$

$$N_5 = 4 \xi_1 \xi_2$$

$$N_6 = 4 \xi_2 \xi_3$$

$$N_7 = 4 \xi_3 \xi_1$$

$$N_8 = 4 \xi_1 \xi_4$$

$$N_9 = 4 \xi_2 \xi_4$$

$$N_{10} = 4 \xi_3 \xi_4$$

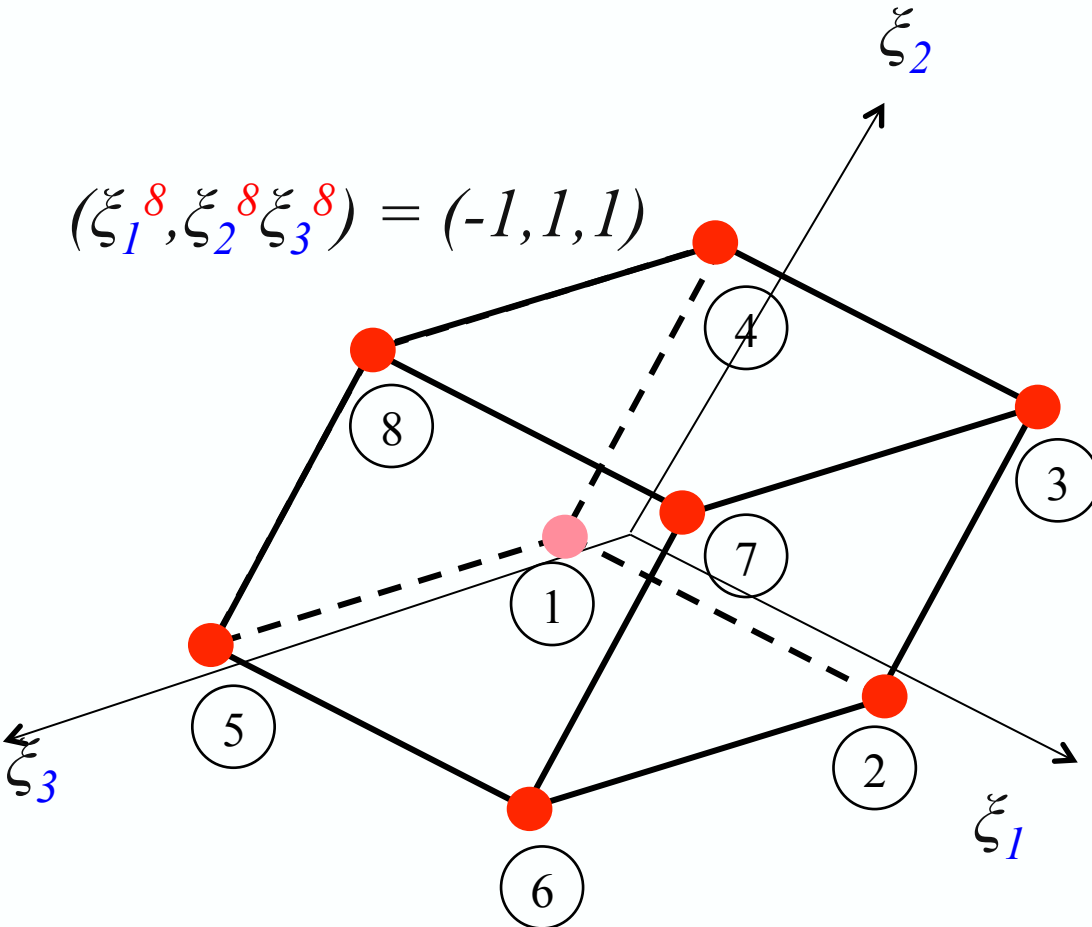
Shape function of a node is product of equations of planes that do not contain that node; e.g. $N_{10} = 4 \xi_3 \xi_4$

Shape functions satisfy: $\sum N_I(\xi) = 1$

8-node isoparametric hexahedron

Shape functions $N_I(\xi)$ ($I = 1, \dots, 8$)

$$(\xi_1^8, \xi_2^8, \xi_3^8) = (-1, 1, 1)$$



$$(\xi_1^6, \xi_2^6, \xi_3^6) = (1, -1, 1)$$

$$N_I(\xi_1, \xi_2, \xi_3) = \frac{1}{8} (1 + \xi_1^I \xi_1) (1 + \xi_2^I \xi_2) (1 + \xi_3^I \xi_3)$$

Shape functions satisfy: $\sum N_I(\xi) = 1$

isoparametric hexahedron

$$\int_{V(\boldsymbol{\xi})} f(\xi_1, \xi_2, \xi_3) d\boldsymbol{\xi} \approx \sum_{I=1}^{n_I} \sum_{J=1}^{n_I} \sum_{K=1}^{n_I} w_I w_J w_K f(\xi_1^I, \xi_2^J, \xi_3^K)$$

1D Gaussian weights
and points of $[-1, 1]$

isoparametric tetrahedron

$$\int_{V(\boldsymbol{\xi})} f(\xi_1, \xi_2, \xi_3, \xi_4) d\boldsymbol{\xi} \approx \sum_{I=1}^{n_I} w_I f(\xi_1^I, \xi_2^I, \xi_3^I, \xi_4^I)$$

3D Gaussian weights
and points calculated in
master element

$n_I = 1$ – accuracy $O(h^2)$

$w_I = 1, \boldsymbol{\xi}^I = (1/4, 1/4, 1/4, 1/4)$

$n_I = 5$ – accuracy $O(h^4)$

$w_1 = -4/5, \boldsymbol{\xi}^1 = (1/4, 1/4, 1/4, 1/4)$

$w_{2,3,4,5} = 9/20, \boldsymbol{\xi}^2 = (1/2, 1/6, 1/6, 1/6)$, others by cyclic symmetry

TORSION OF A BAR IN 3D; DISPLACEMENT APPROACH

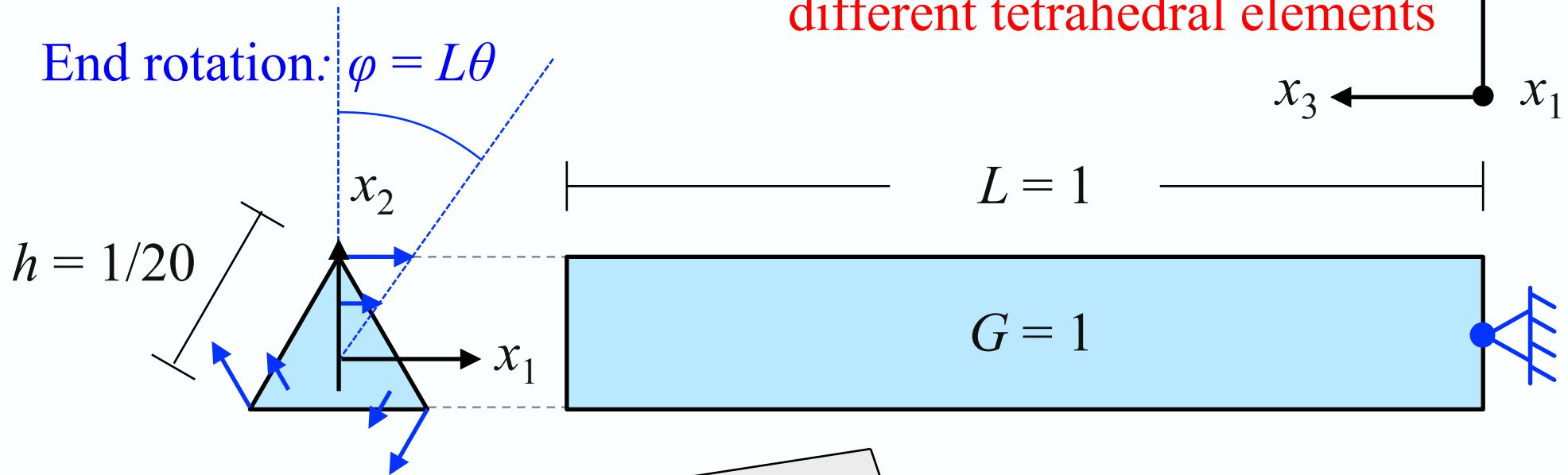
- There are problems in elasticity that are **more easily solved with a stress-based formulation** (and not with the displacement-based formulation that we have seen so far).
- One such case is the **torsion of a prismatic bar with arbitrary section; away from the two ends** (where a boundary layer may develop) **the stresses are independent of the axial coordinate** (depend solely on the two in-plane coordinates); the **displacements depend on all three coordinates**
- Here we solve the torsion problem of a triangular section bar **using a displacement-based FEM formulation and 3D tetrahedral elements**; we **compare the solution to the analytical one** (obtained using the stress-based formulation) and also **show the development of a boundary layer** at one end in the case where warping is prevented (full clamping).

End rotation at $x_3=L$ imposed through the essential boundary conditions:

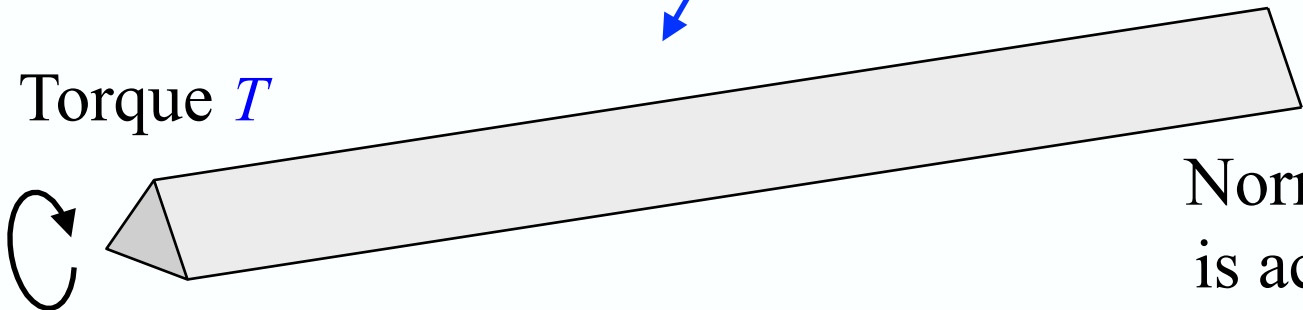
$$u_1(x_1, x_2, L) = x_2 \varphi, \quad u_2(x_1, x_2, L) = -x_1 \varphi$$

Prismatic bar analyzed using an FEM **displacement-based** discretization based on **different tetrahedral elements**

End rotation: $\varphi = L\theta$

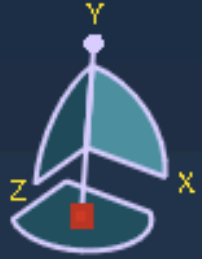
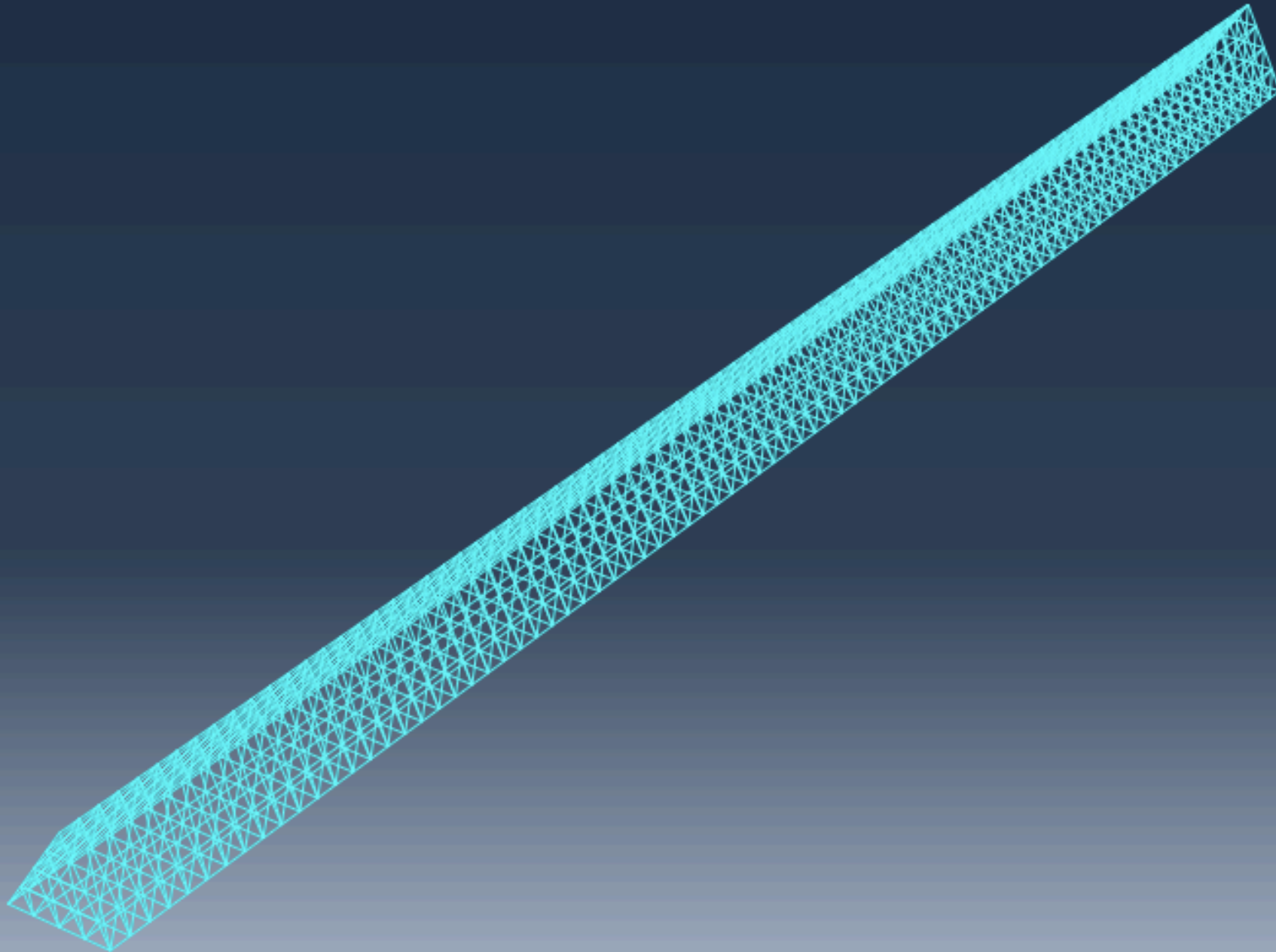


Torque T

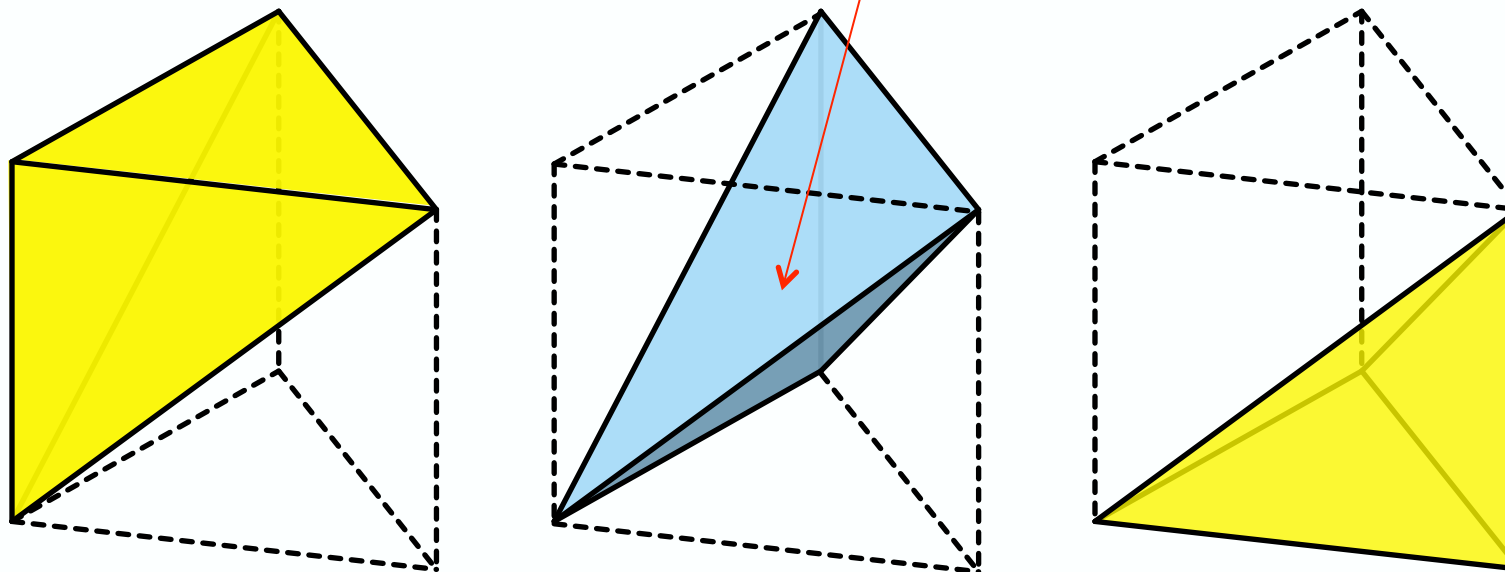


Normal stress-free ($\sigma_{33}=0$) at $x_3=0$ is accounted through the essential boundary conditions:

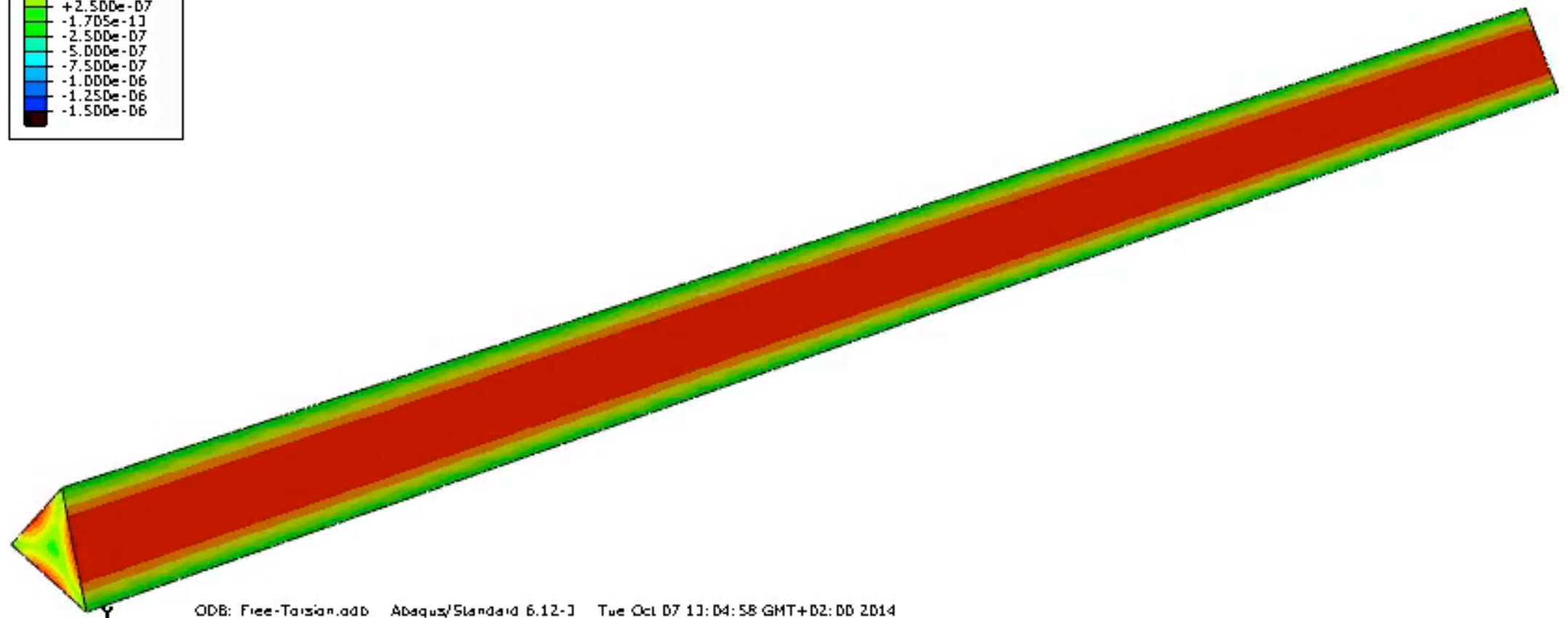
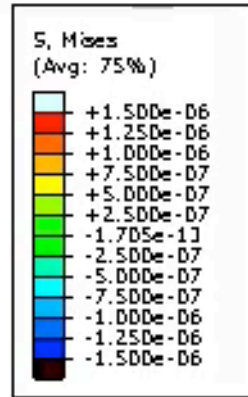
Prismatic bar subjected to **rate of twist θ** $u_1(x_1, x_2, 0) = u_2(x_1, x_2, 0) = 0$



NOTE on meshing: a straightforward division of the pentahedral prism in three tetrahedra gives one **ill-conditioned element**; as a result the code automatically changes meshes to have well-behaved elements (with gradients of the same order in all directions). This explains the rather bizzare number of total elements in the 3D prism



Scale Factor: +1.00



ODB: Free-Torsion.odb Abaqus/Standard 6.12-1 Tue Oct 07 11:04:58 GMT+02:00 2014

Step: Step-1
 Increment 1: Step Time = 1.000
 Primary Var: S, Mises
 Deformed Var: U Deformation Scale Factor: +1.000e+03

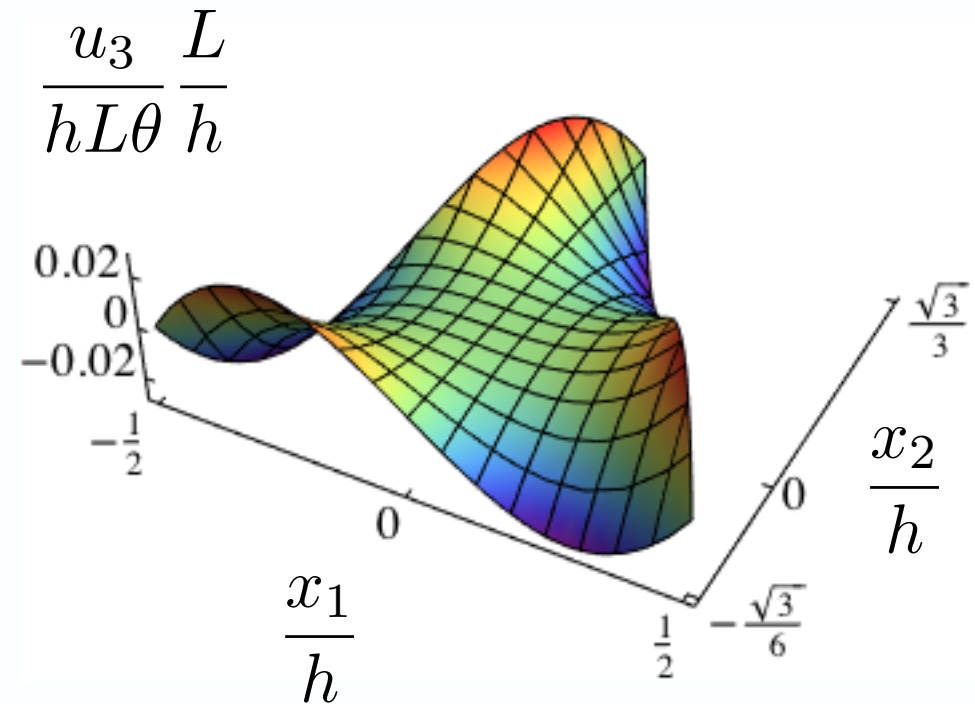
(non-dimensionalized) **warping** (i.e. $u_3(x_1, x_2)$) of each cross-section

Global displacement field (analytical)

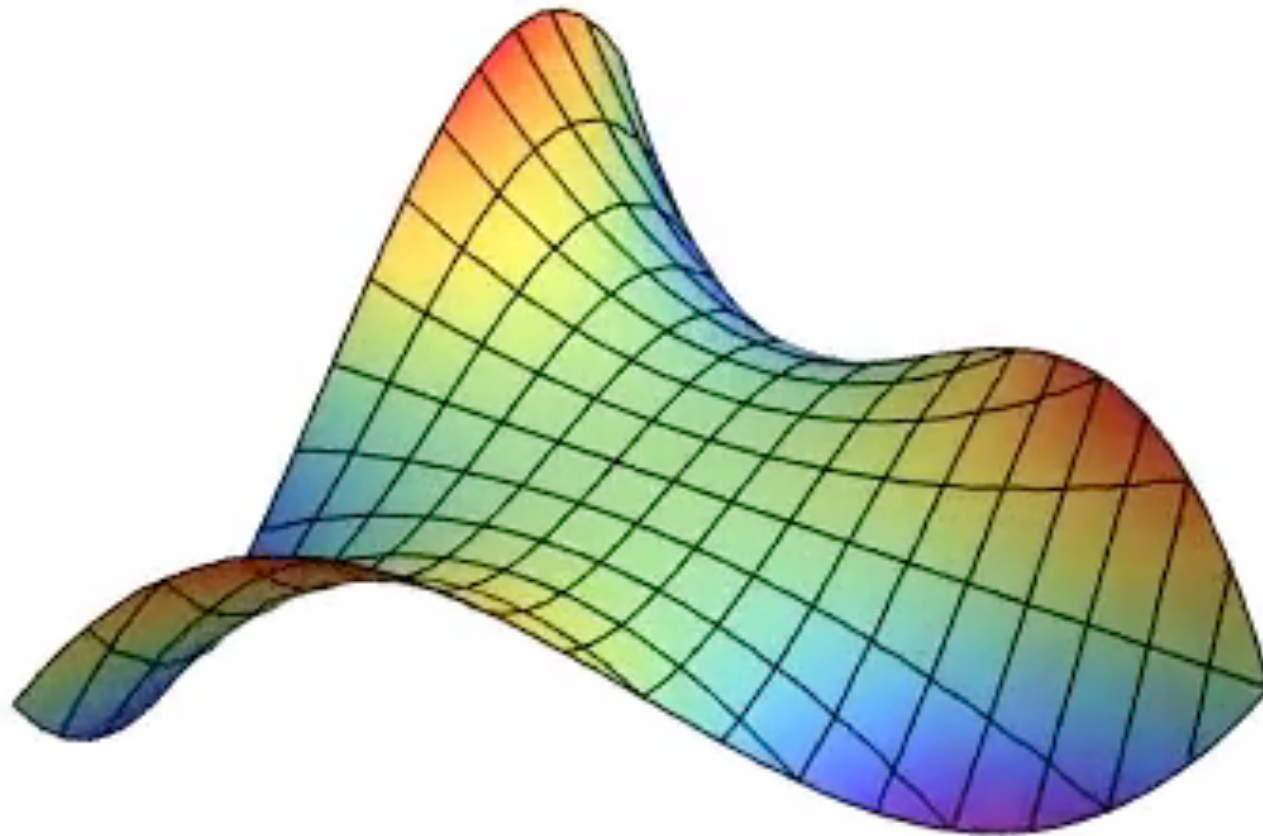
$$\frac{u_1}{hL\theta} = \frac{x_3}{L} \frac{x_2}{h}$$

$$\frac{u_2}{hL\theta} = -\frac{x_3}{L} \frac{x_1}{h}$$

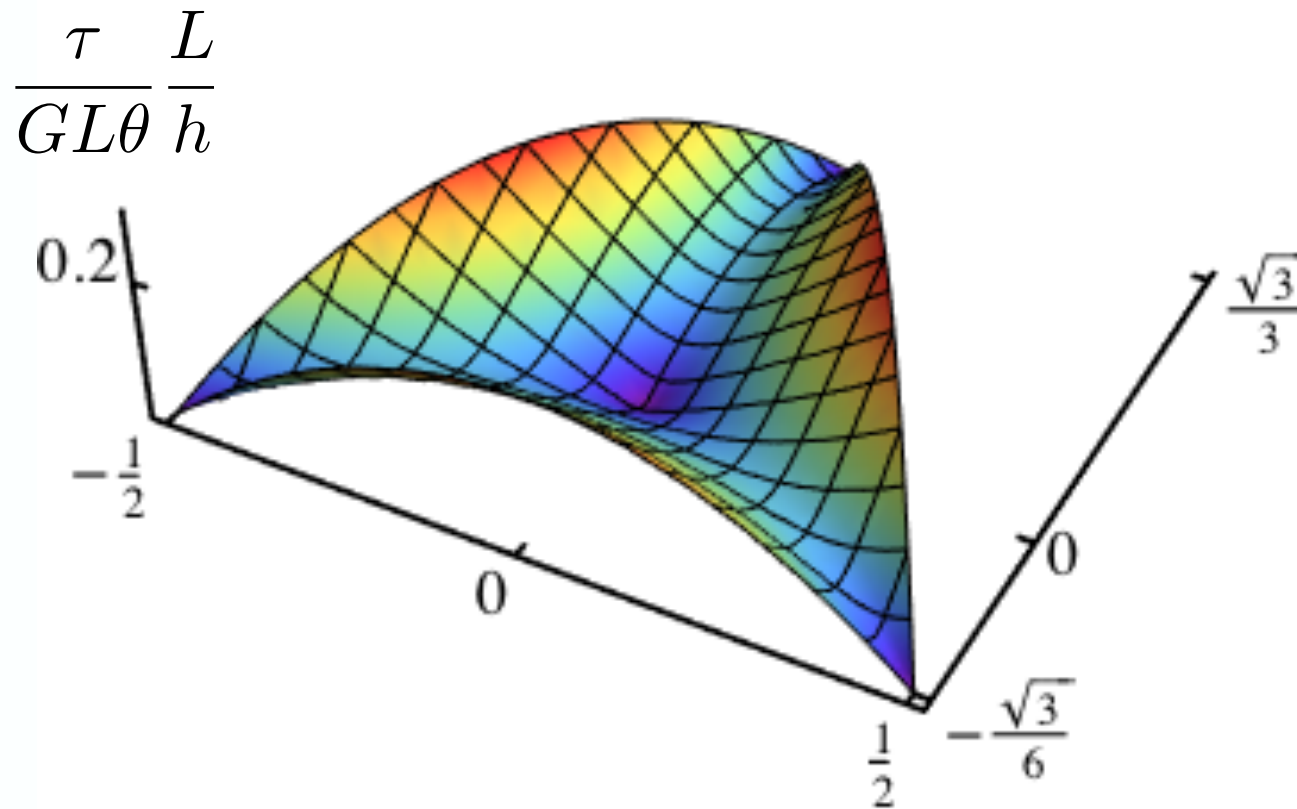
$$\frac{u_3}{hL\theta} = \sqrt{3} \frac{h}{L} \left[\frac{1}{3} \left(\frac{x_1}{h} \right)^3 - \frac{x_1}{h} \left(\frac{x_2}{h} \right)^2 \right]$$



Notice the C_3 symmetry of the warping (i.e. $u_3(x_1, x_2)$) distribution;
 Warping is zero at middle of the section and of the edges

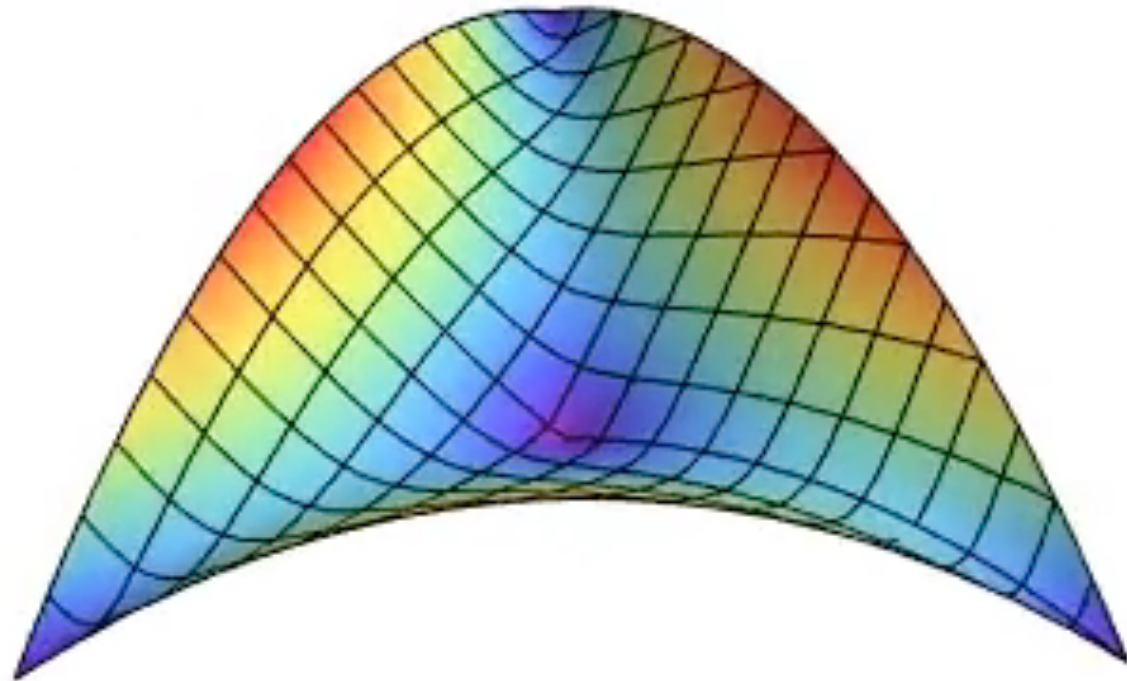


(non-dimensionalized) shear stress $\tau = [(\sigma_{31})^2 + (\sigma_{32})^2]^{1/2}$ on the cross-section



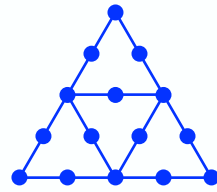
Notice the C_3 symmetry of the stress distribution;
 Stress is zero at middle and at corners and maximum at middle of edges

STRESS STATE AT EACH SECTION (ONLY SHEAR EXISTS)

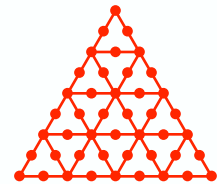


analytical solution $\frac{\tau}{GL\theta} \frac{L}{\theta} \Big|_{x_1=0} = \frac{\sqrt{3}}{2} \left| \frac{x_2}{h} \left(\frac{x_2}{h} - \frac{\sqrt{3}}{3} \right) \right|$

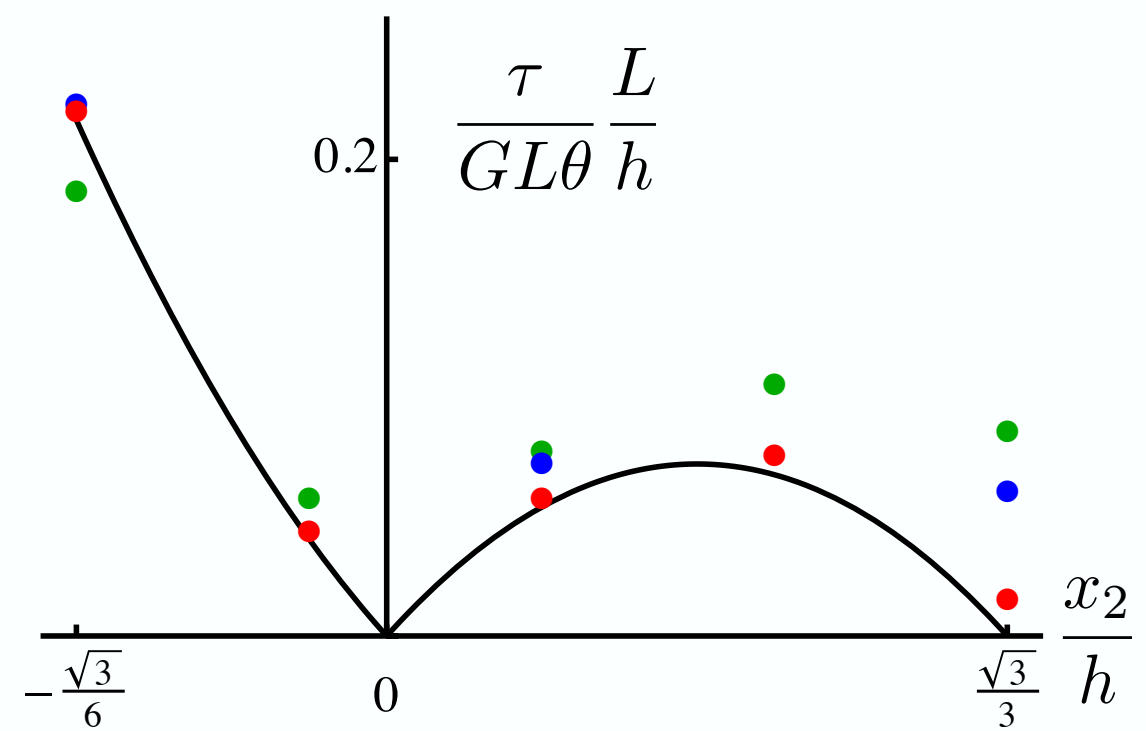
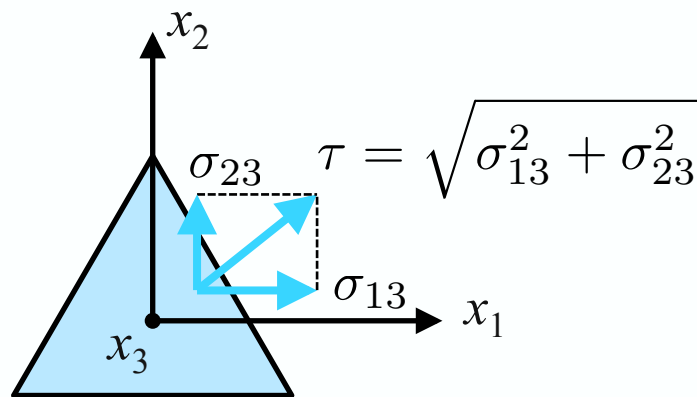
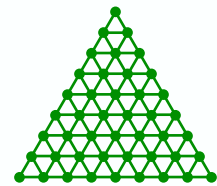
2 layers of quadratic elements



4 layers of quadratic elements



8 layers of linear elements

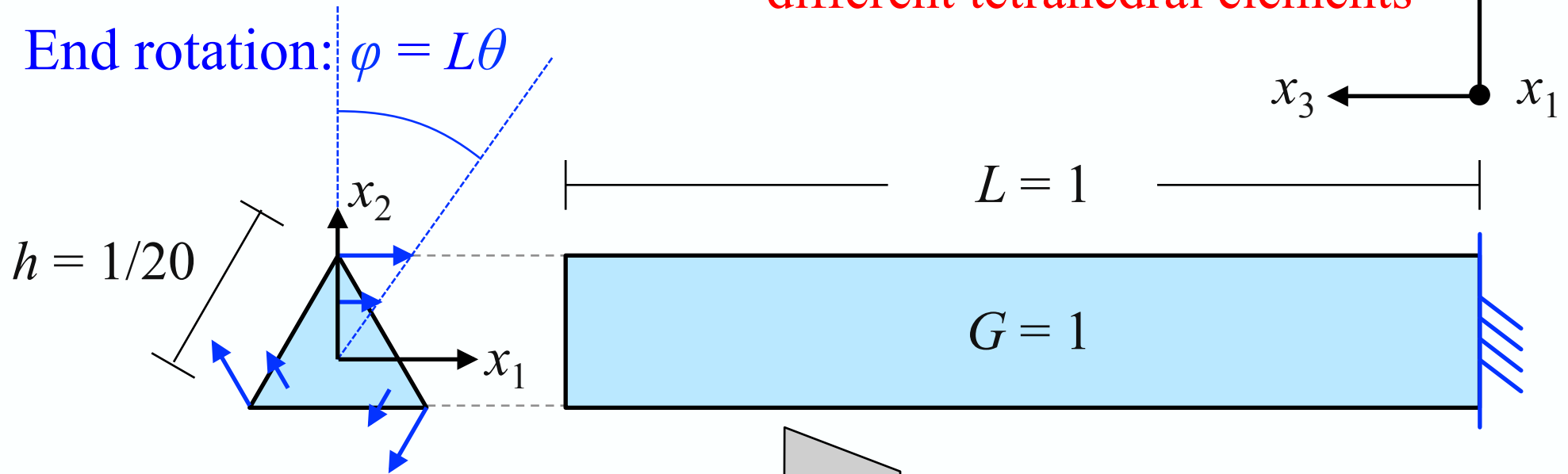


End rotation at $x_3=L$ imposed through the essential boundary conditions:

$$u_1(x_1, x_2, L) = x_2 \varphi, \quad u_2(x_1, x_2, L) = -x_1 \varphi$$

Prismatic bar analyzed using an FEM **displacement-based** discretization based on **different tetrahedral elements**

End rotation: $\varphi = L\theta$



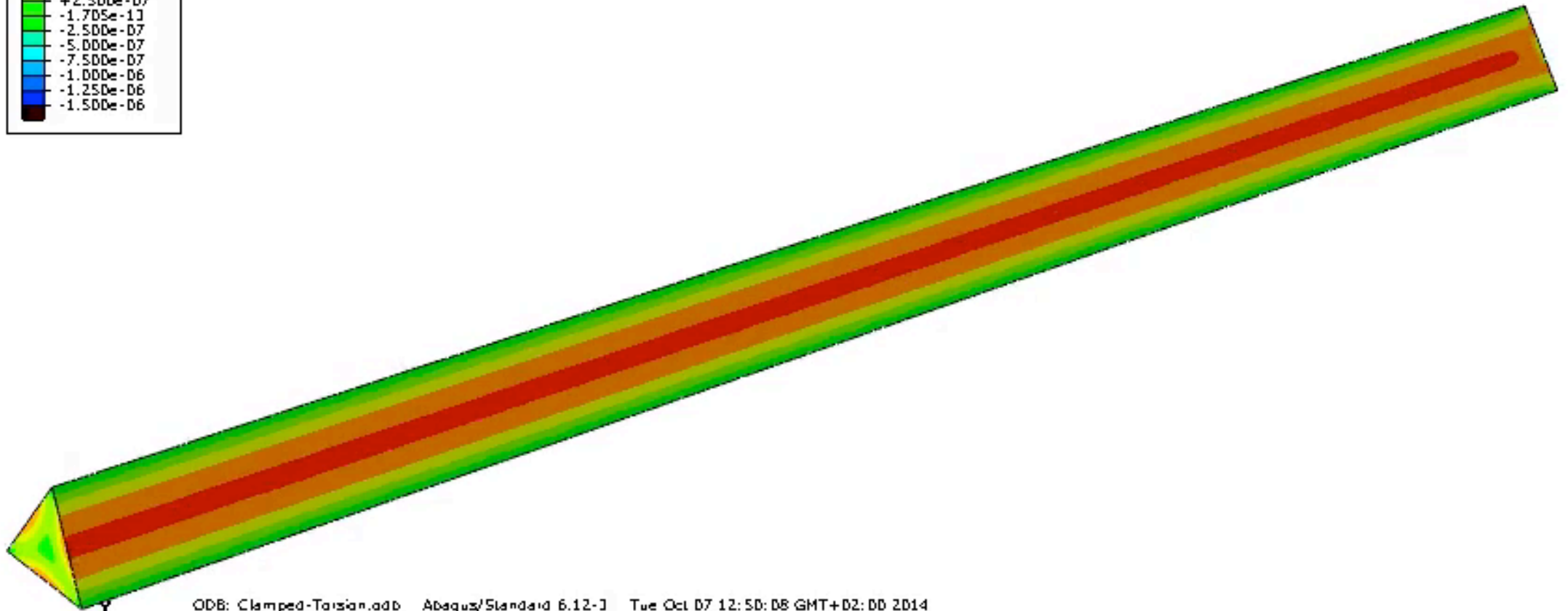
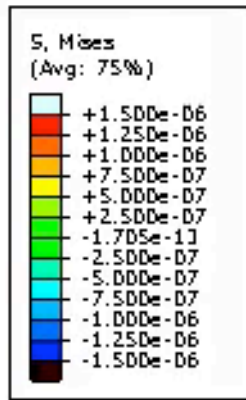
Torque T

Clamped support at $x_3=0$:
 $u_i(x_1, x_2, 0) = 0 \quad (i = 1, 2, 3),$

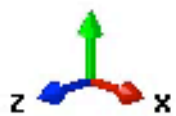
Prismatic bar subjected to **rate of twist θ**

Boundary layer develops!

Scale Factor: +0.72



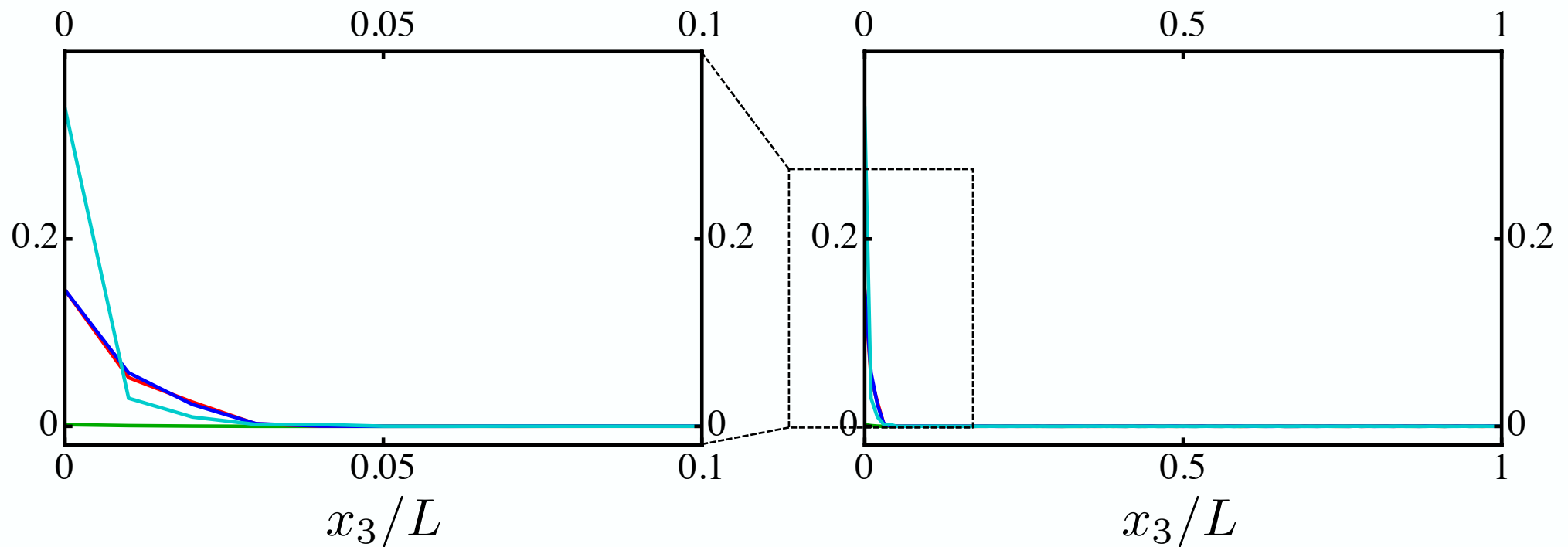
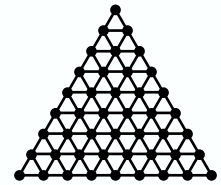
ODB: Clamped-Torsion.odb Abaqus/Standard 6.12-3 Tue Oct 07 12:50:08 GMT+02:00 2014



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 Increment 1: Step Time = 1.000
 Primary Var: S, Mises
 Deformed Var: U Deformation Scale Factor: +1.000e+03

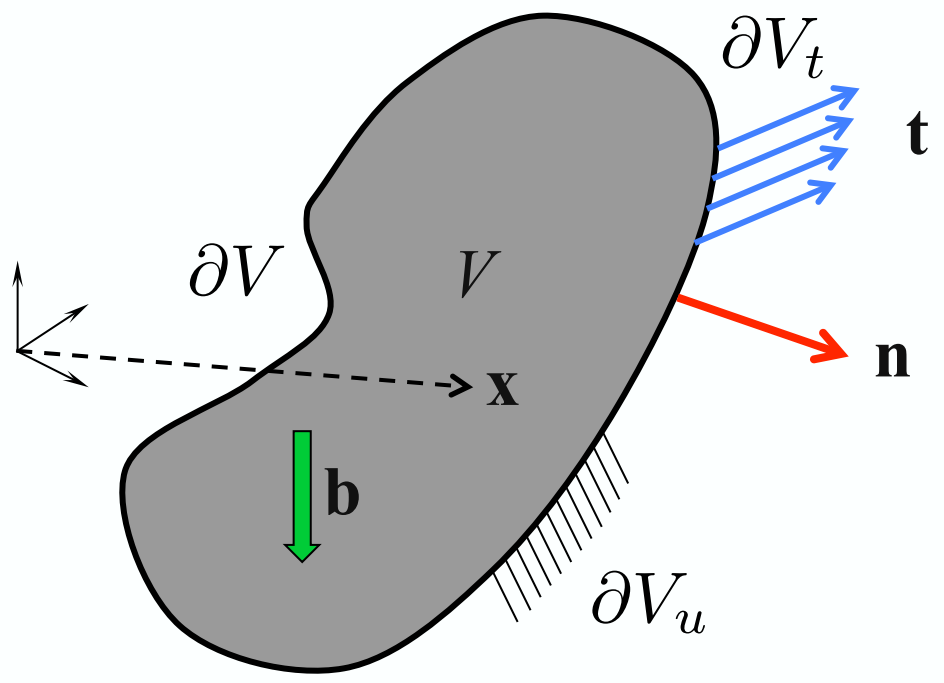
$$\frac{\sigma_{11}}{GL\theta} \frac{L}{h} \quad \frac{\sigma_{22}}{GL\theta} \frac{L}{h} \quad \frac{\sigma_{12}}{GL\theta} \frac{L}{h} \quad \frac{\sigma_{33}}{GL\theta} \frac{L}{h}$$

8 layers of linear elements



Since warping is prevented at $x_3=0$, normal (and a tiny in-plane shear) stresses develop in a **boundary layer** of the order of the section thickness

PRINCIPLE OF MINIMUM COMPLEMENTARY ENERGY



Solid occupies domain: V

Domain boundary: ∂V

Body forces: \mathbf{b}

Surface traction: \mathbf{t}

Surface normal (outward): \mathbf{n}

Traction prescribed on: ∂V_t

Displacement prescribed on: ∂V_u

Position vector: \mathbf{x}

Complementary energy density: $W^*(\boldsymbol{\sigma})$

Strain-stress:
$$\epsilon_{ij} = \frac{\partial W^*}{\partial \sigma_{ij}}$$

(general nonlinear elastic material)

Complementary : $\mathcal{P}^* = \mathcal{P}_{int}^* + \mathcal{P}_{ext}^*$

Internal : $\mathcal{P}_{int}^* = \int_V W^*(\sigma_{ij}) dV ; \quad \epsilon_{ij} = \frac{\partial W^*}{\partial \sigma_{ij}}$

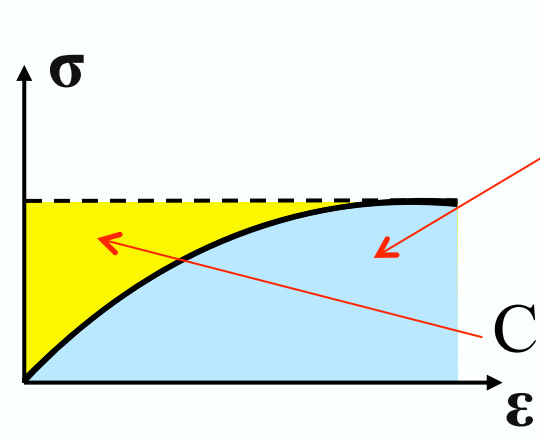
External : $\mathcal{P}_{ext}^* = - \int_V b_i u_i dV - \int_{\partial V_u} t_i u_i dS$

$\mathcal{P}^*(\boldsymbol{\sigma} + \epsilon \delta \boldsymbol{\sigma}) \geq \mathcal{P}^*(\boldsymbol{\sigma}) ;$
 Equilibrium at minimum over all
statically admissible stress fields

$\nabla \bullet \delta \boldsymbol{\sigma}(\mathbf{x}) = 0, \quad \forall \mathbf{x} \in V$
 $\mathbf{n} \bullet \delta \boldsymbol{\sigma}(\mathbf{x}) = 0, \quad \forall \mathbf{x} \in \partial V_t$

$\frac{d}{d\epsilon} [\mathcal{P}^*(\boldsymbol{\sigma} + \epsilon \delta \boldsymbol{\sigma})]_{\epsilon=0} = 0 ; \quad \text{extremum (1)}$

$\frac{d^2}{d\epsilon^2} [\mathcal{P}^*(\boldsymbol{\sigma} + \epsilon \delta \boldsymbol{\sigma})]_{\epsilon=0} > 0 ; \quad \text{minimum (2)}$



Potential energy: $W(\epsilon) = \int_0^\epsilon \sigma_{ij} d\epsilon_{ij}$

Complementary energy: $W^*(\sigma) = \int_0^\sigma \epsilon_{ij} d\sigma_{ij}$

Equilibrium : $\frac{\partial \sigma_{ij}}{\partial x_i} + b_j = 0, \mathbf{x} \in V$

Boundary : $t_j = n_i \sigma_{ij}, \mathbf{x} \in \partial V_t$

static
admissibility
condition

Constitutive (Linear) : $\epsilon_{ij} = L_{ijkl}^{-1} \sigma_{kl}$

Complementary (Linear) : $W^* = \int_0^\sigma \epsilon_{ij} d\sigma_{ij} = \frac{1}{2} L_{ijkl}^{-1} \sigma_{ij} \sigma_{kl}$

TORSION OF A BAR IN 3D; STRESS-BASED APPROACH

$$\begin{aligned} \mathcal{P}_{int}^* &= \int_V W^*(\sigma_{ij}) dV = \int_V \frac{1}{2G} \left[(\sigma_{32})^2 + (\sigma_{31})^2 \right] dV \\ &= L \int_A \frac{1}{2G} \left[\left(\frac{\partial \psi}{\partial x_1} \right)^2 + \left(\frac{\partial \psi}{\partial x_2} \right)^2 \right] dA \end{aligned}$$

$$\mathcal{P}_{ext}^* = - \int_{\partial V_u} t_i u_i dS = - \int_A [\sigma_{31} u_1 + \sigma_{32} u_2] dA$$

$$= - \int_A [\sigma_{31}(\phi x_2) + \sigma_{32}(-\phi x_1)] dA = -T\phi = -TL\theta$$

$$\text{Torque } T = \int_A \left[\left(\frac{\partial \psi}{\partial x_2} \right) x_2 + \left(\frac{\partial \psi}{\partial x_1} \right) x_1 \right] dA = 2 \int_A \psi dA$$

Torsion for bar of arbitrary section A using (Prandtl) stress potential ψ :

$$\mathcal{P}_{int}^* = L \int_A \frac{1}{2G} \left[\left(\frac{\partial \psi}{\partial x_1} \right)^2 + \left(\frac{\partial \psi}{\partial x_2} \right)^2 \right] dA$$

$$\mathcal{P}_{ext}^* = -2L\theta \int_A \psi dA$$

Admissibility : $\psi = 0 \quad \forall \mathbf{x} \in \partial A \quad (\sigma_{31}n_1 + \sigma_{32}n_2 = 0)$

Euler – Lagrange : $\frac{\partial^2 \psi}{\partial x_1 \partial x_1} + \frac{\partial^2 \psi}{\partial x_2 \partial x_2} - 2G\theta = 0 \quad \forall \mathbf{x} \in A$

Prandtl function $\psi(x_1, x_2) = \frac{\sqrt{3}\theta G}{12h} \left(\sqrt{3}h (x_1^2 + x_2^2) + 6x_1^2x_2 - 2x_2^3 \right)$

$$\sigma_{31}(x_1, x_2) = \frac{\partial \psi}{\partial x_2} = Gh\theta \left[\frac{\sqrt{3}}{2} \left(\left(\frac{x_1}{h} \right)^2 - \left(\frac{x_2}{h} \right)^2 \right) + \frac{x_2}{2h} \right]$$

$$\sigma_{32}(x_1, x_2) = -\frac{\partial \psi}{\partial x_1} = -Gh\theta \left(\frac{x_1}{2h} + \sqrt{3} \frac{x_1x_2}{h^2} \right)$$

Shear stress $\tau(x_1, x_2) \equiv \sqrt{\sigma_{31}^2 + \sigma_{32}^2}$

NOTE : $\sigma_{11} = \sigma_{22} = \sigma_{12} = \sigma_{33} = 0$

FEM DISCRETIZATION USING COMPLEMENTARY ENERGY

We discretize the (Prandtl) stress potential $\psi(x_1, x_2)$ over section A

$$\mathcal{P}^* = L \int_A \left\{ \frac{1}{2G} \left[\left(\frac{\partial \psi}{\partial x_1} \right)^2 + \left(\frac{\partial \psi}{\partial x_2} \right)^2 \right] - 2\theta\psi \right\} dA$$

$$\frac{1}{L} \delta \mathcal{P}^* = \int_A \left\{ \frac{1}{G} \left[\frac{\partial \psi}{\partial x_1} \frac{\partial \delta \psi}{\partial x_1} + \frac{\partial \psi}{\partial x_2} \frac{\partial \delta \psi}{\partial x_2} \right] - 2\theta \delta \psi \right\} dA = 0$$

discretization : $\psi = \mathbf{N} \mathbf{q}_e, \quad \mathbf{q}_e = [\Psi^1, \Psi^2, \Psi^3, \dots]^T$

discretization : $\left[\frac{\partial \psi}{\partial x_1}, \frac{\partial \psi}{\partial x_2} \right]^T = \mathbf{B} \mathbf{q}_e$

element : $\mathbf{k}_e = \frac{1}{G} \int_{A_e} [\mathbf{B}^T \mathbf{B}] dA, \quad \mathbf{f}_e = 2\theta \int_{A_e} [\mathbf{N}^T] dA$

PROGRAM CHANGE:

NOV 9: Lecture 6 at 10:45 in PC64

NOV 9: Lecture 7 at 16:15 IN PC65

NOV 16: PC 6 at 10:45 in Salle Info n°31

NOV 16: PC 7 at 16:15 in Salle Info n°31