

TOPICS COVERED IN THIS LECTURE

1. BAR MODEL (1D) – LINEAR (**REVIEW**) AND HIGHER ORDER ELEMENTS
2. BAR MODEL (1D) – MASTER ELEMENT & ISOPARAMETRIC INTERPOLATION
3. PLANAR TRUSSES (2D)
4. SPACE FRAMES (3D)
5. THERMAL LOADING OF TRUSSES

Potential : $\mathcal{P} = \mathcal{P}_{int} + \mathcal{P}_{ext}$

Internal : $\mathcal{P}_{int} = \int_V \frac{1}{2} \sigma(x) \epsilon(x) dV = \frac{1}{2} \int_0^L E A(x) \left(\frac{du}{dx} \right)^2 dx$

External : $\mathcal{P}_{ext} = - \int_V \rho g u(x) dV = - \int_0^L \rho g A(x) u(x) dx$

Numerical solution of equilibrium problem – **Raleigh-Ritz method**:

Instead of **minimizing energy in an infinite dimensional space**, we do **minimize in a finite dimensional space**, in which case we end up with an **algebraic problem**. Since energy is **quadratic in $u(x)$** , system is **linear!**

$$\mathcal{P}(u^{app}(x)) = \mathcal{P}(\mathbf{Q}); \quad u^{app} = \sum_{i=1}^{i=n} Q_i N_i(x), \quad \mathbf{Q} \equiv [Q_1, Q_2, \dots, Q_n]$$

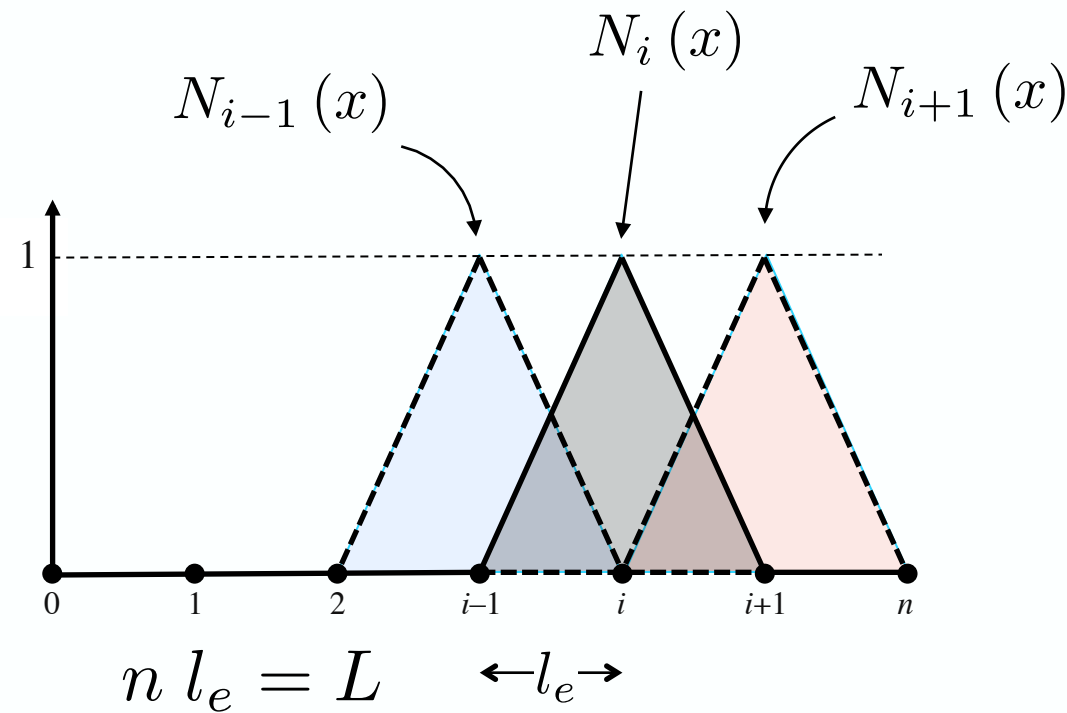
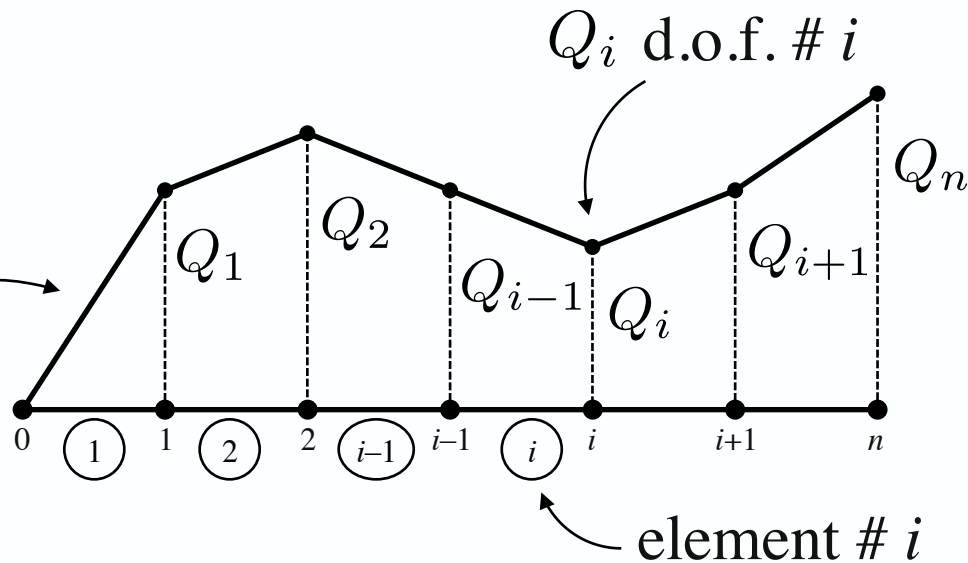
$$\frac{\partial \mathcal{P}(\mathbf{Q})}{\partial Q_i} = \int_0^L \left[EA \frac{du^{app}}{dx} \frac{\partial}{\partial Q_i} \left(\frac{du^{app}}{dx} \right) - \rho g A \frac{\partial u^{app}}{\partial Q_i} \right] dx = 0$$

$$\sum_{j=1}^{j=n} \left\{ \int_0^L \left[EA \frac{dN_i}{dx} \frac{dN_j}{dx} \right] dx \right\} Q_j - \int_0^L [\rho g A N_i] dx = 0$$

$$\sum_{j=1}^{j=n} K_{ij} Q_j - F_i = 0 ; \quad (\text{in compact form : } \mathbf{KQ} = \mathbf{F})$$

$$K_{ij} \equiv \int_0^L \left[EA \frac{dN_i}{dx} \frac{dN_j}{dx} \right] dx , \quad F_i \equiv \int_0^L [\rho g A N_i] dx$$

Stiffness matrix: \mathbf{K} , Force vector: \mathbf{F} , Degrees of Freedom: \mathbf{Q}



$$u^{\text{app}}(x) = \sum_i Q_i N_i(x)$$

Shape function $N_i(x)$

Easy physical interpretation of d.o.f. (degree of freedom) Q_i at **node** x_i : due to its construction, $u^{\text{app}}(x_i) = Q_i$

Shape functions $N_i(x)$ have **compact support**: $N_i(x_i) = 1, N_i(x_{i-1}) = N_i(x_{i+1}) = 0$.
Compactness of support of shape function **great advantage of FEM**

$$\begin{bmatrix}
 K_{11} & K_{12} & 0 & 0 & 0 \\
 K_{21} & K_{22} & K_{23} & 0 & 0 \\
 0 & K_{32} & K_{33} & K_{34} & 0 \\
 0 & 0 & K_{43} & K_{44} & K_{45} \\
 0 & 0 & 0 & K_{54} & K_{55}
 \end{bmatrix}$$

$$K_{ii} = \int_{x_{i-1}}^{x_{i+1}} E A(x) \left(\frac{1}{l_e} \right)^2 dx$$

$$K_{i \ i+1} = - \int_{x_i}^{x_{i+1}} E A(x) \left(\frac{1}{l_e} \right)^2 dx$$

Stiffness matrix **K** is **banded**, i.e. populated about the diagonal. Note that the F.E.M. (compact support) shape functions $N_i(x)$ **do not have to be piecewise linear**, as calculated above; a wide selection is possible.

We will show here how we can generalize to **higher order polynomials**.

$$\mathcal{P}(\mathbf{Q}) = \mathcal{P}_{int}(\mathbf{Q}) + \mathcal{P}_{ext}(\mathbf{Q})$$

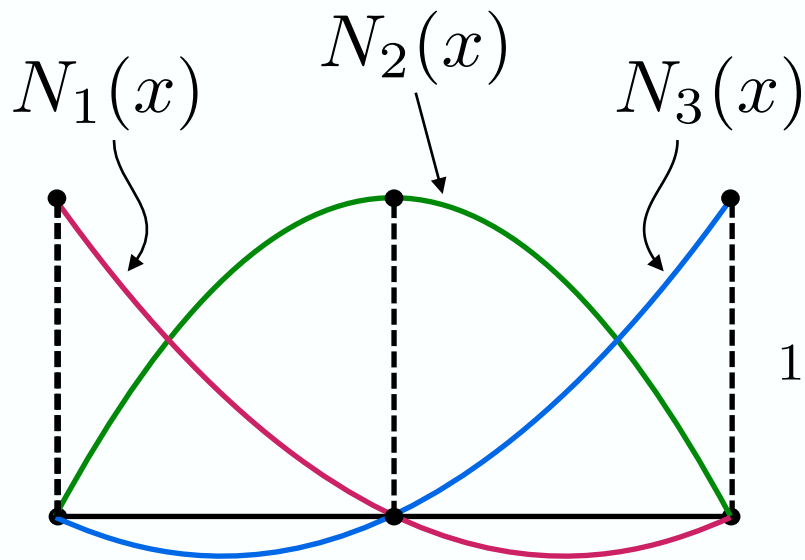
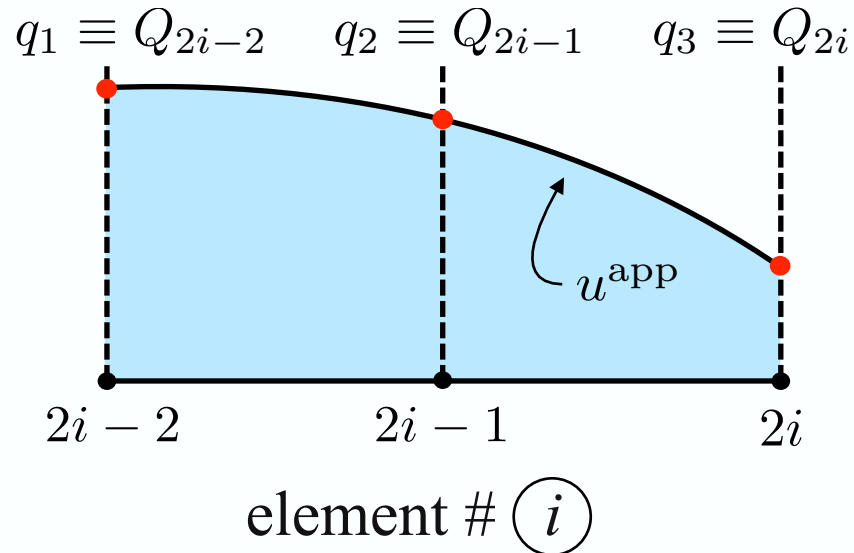
$$\mathcal{P}_{int}(\mathbf{Q}) = \sum_e \mathcal{P}_{int}^e; \quad \mathcal{P}_{int}^e = \int_{l_e} \left[\frac{1}{2} E A(x) \left(\sum_{i=1}^{i=m} q_i \frac{dN_i}{dx} \right)^2 \right] dx = \frac{1}{2} \mathbf{q}_e^T \mathbf{k}_e \mathbf{q}_e$$

$$\mathcal{P}_{ext}(\mathbf{Q}) = \sum_e \mathcal{P}_{ext}^e; \quad \mathcal{P}_{ext}^e = - \int_{l_e} \left[\rho g A(x) \left(\sum_{i=1}^{i=m} q_i N_i(x) \right) \right] dx = -\mathbf{q}_e^T \mathbf{f}_e$$

$$[\mathbf{k}_e]_{ij} \equiv \int_{l_e} \left[E A(x) \left(\frac{dN_i}{dx} \right) \left(\frac{dN_j}{dx} \right) \right] dx; \quad \mathbf{k}_e \text{ element stiffness } (m \times m)$$

$$[\mathbf{f}_e]_i \equiv - \int_{l_e} [\rho g A(x) N_i(x)] dx; \quad \mathbf{f}_e \text{ element force } (m)$$

Element stiffness \mathbf{k}_e ($m \times m$) & force \mathbf{f}_e (m) for $(m-1)$ polynomial interpolation

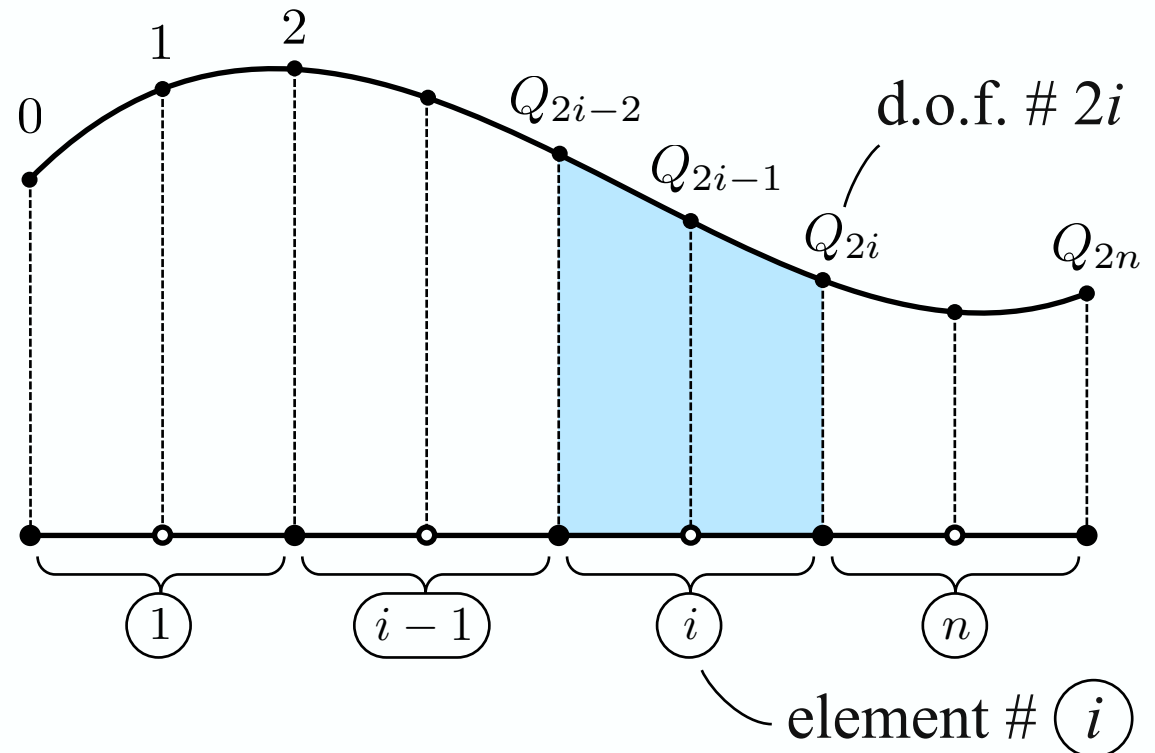


Case $m=3$ (2nd order polynomial interp.)

$$u^{app}(x) = q_1 N_1(x) + q_2 N_2(x) + q_3 N_3(x)$$

Local degree of freedom $\mathbf{q}^T_e = [q_1, q_2, q_3]$

In each element: $u^{app}(x_i) = q_i$ ($i=1,2,3$)



Case of a **quadratic** polynomial interpolation ($m=3$)

Recall requirement: $u^{app}(x_i) = q_1 N_1(x_i) + q_2 N_2(x_i) + q_3 N_3(x_i) = q_i$

$$N_1(x_1) = 1, \quad N_1(x_2) = 0, \quad N_1(x_3) = 0$$

$$N_1(x) = [(x - x_2)(x - x_3) / [(x_1 - x_2)(x_1 - x_3)]]$$

$$N_2(x_1) = 0, \quad N_2(x_2) = 1, \quad N_2(x_3) = 0$$

$$N_2(x) = [(x - x_1)(x - x_3) / [(x_2 - x_1)(x_2 - x_3)]]$$

$$N_3(x_1) = 0, \quad N_3(x_2) = 0, \quad N_3(x_3) = 1$$

$$N_3(x) = [(x - x_1)(x - x_2) / [(x_3 - x_1)(x_3 - x_2)]]$$

You see now the easy extension to higher order interpolation!

Case of a **cubic** polynomial interpolation ($m=4$)

$$N_1(x_1) = 1, \quad N_1(x_2) = 0, \quad N_1(x_3) = 0, \quad N_1(x_4) = 0$$

$$N_1(x) = [(x - x_2)(x - x_3)(x - x_4)] / [(x_1 - x_2)(x_1 - x_3)(x_1 - x_4)]$$

$$N_2(x_1) = 0, \quad N_2(x_2) = 1, \quad N_2(x_3) = 0, \quad N_2(x_4) = 0$$

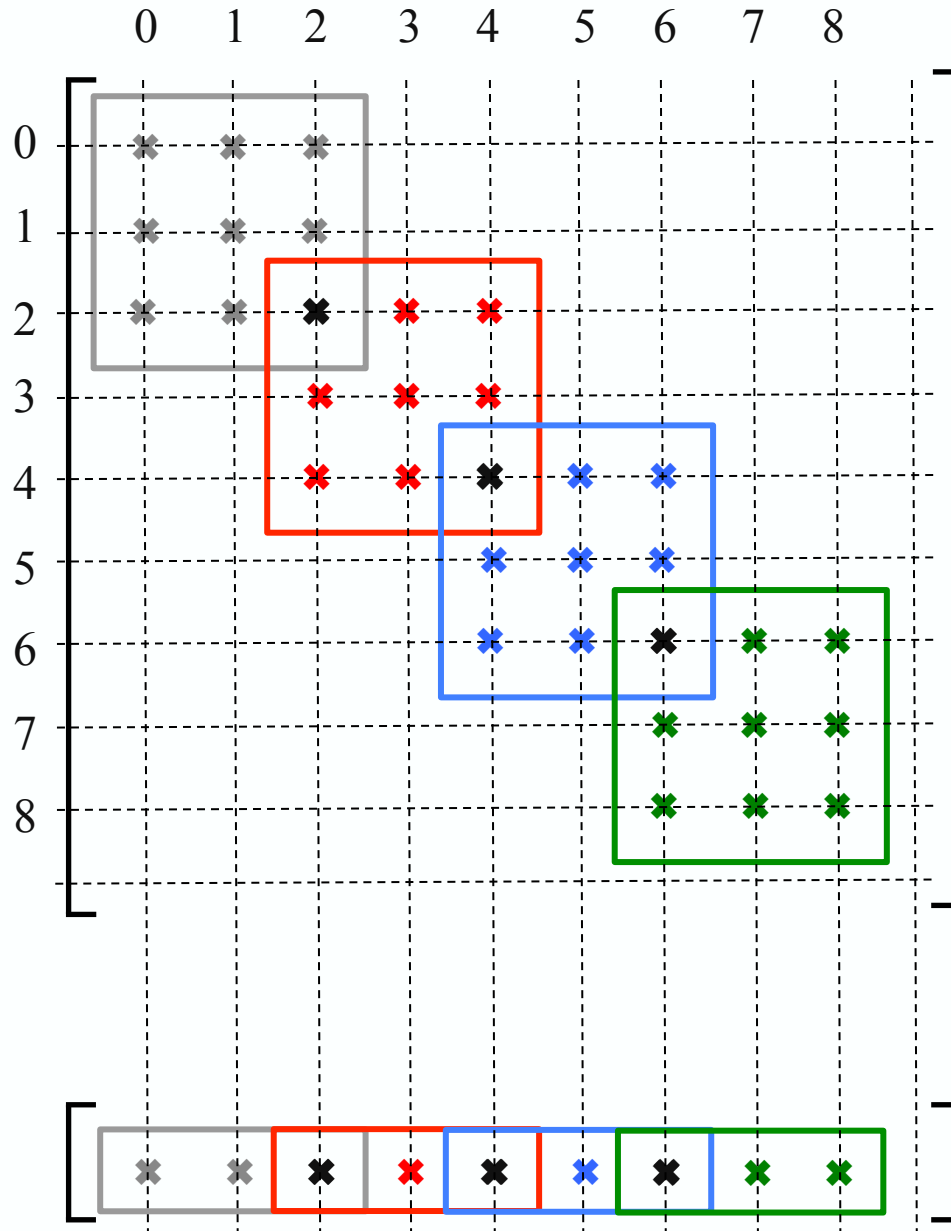
$$N_2(x) = [(x - x_1)(x - x_3)(x - x_4)] / [(x_2 - x_1)(x_2 - x_3)(x_2 - x_4)]$$

$$N_3(x_1) = 0, \quad N_3(x_2) = 0, \quad N_3(x_3) = 1, \quad N_3(x_4) = 0$$

$$N_3(x) = [(x - x_1)(x - x_2)(x - x_4)] / [(x_3 - x_1)(x_3 - x_2)(x_3 - x_4)]$$

$$N_4(x_1) = 0, \quad N_4(x_2) = 0, \quad N_4(x_3) = 0, \quad N_4(x_4) = 1$$

$$N_4(x) = [(x - x_1)(x - x_2)(x - x_3)] / [(x_4 - x_1)(x_4 - x_2)(x_4 - x_3)]$$



Assembling global stiffness matrix K and global force vector F from element stiffness matrix k_e and element force vector f_e : add local components in the appropriate place of global counterparts (for quadratic polynomial interpolation, $m=3$)

local no. \rightarrow global no.

1 \rightarrow 0, 2 \rightarrow 1, 3 \rightarrow 2 for element 1

1 \rightarrow 2, 2 \rightarrow 3, 3 \rightarrow 4 for element 2

1 \rightarrow 4, 2 \rightarrow 5, 3 \rightarrow 6 for element 3

1 \rightarrow 6, 2 \rightarrow 7, 3 \rightarrow 8 for element 4

It is very convenient to write shape functions in a **master element** with respect to a **normalized coordinate** (ξ)

$$\begin{array}{ccc}
 i = 1 & & i = 2 \\
 \bullet & \text{---} & \bullet \\
 x = 0 & & x = 1
 \end{array}
 \quad m = 2 ; \quad \begin{array}{l} N_1(\xi) = 1 - \xi \\ N_2(\xi) = \xi \end{array}$$

$$\begin{array}{ccccc}
 i = 1 & & i = 2 & & i = 3 \\
 \bullet & \text{---} & \bullet & \text{---} & \bullet \\
 x = -1 & & x = 0 & & x = 1
 \end{array}
 \quad m = 3 ; \quad \begin{array}{l} N_1(\xi) = \xi(\xi - 1)/2 \\ N_2(\xi) = (1 - \xi)(\xi + 1)/2 \\ N_3(\xi) = \xi(\xi + 1)/2 \end{array}$$

$$[\mathbf{k}_e]_{ij} \equiv \int_{\xi} \left[E A(x(\xi)) (dN_i/d\xi) (dN_j/d\xi) (dx/d\xi)^{-1} \right] d\xi$$

$$[\mathbf{f}_e]_i \equiv - \int_{\xi} \left[\rho g A(x(\xi)) N_i(\xi) (dx/d\xi)^{-1} \right] d\xi$$

Question: what do we choose for $x(\xi)$?

Answer: (easy) same representation as for displacement!

This type of parametrization that uses the **same** interpolation scheme for both the **displacement** and the **geometric coordinates** is called **isoparametric** representation and is widely used in F.E.M.

$$u(\xi) = \sum_{i=1}^{i=m} u_i N_i(\xi) , \quad u(\xi_i) = u_i ; \quad \text{d.o.f. at node } i$$

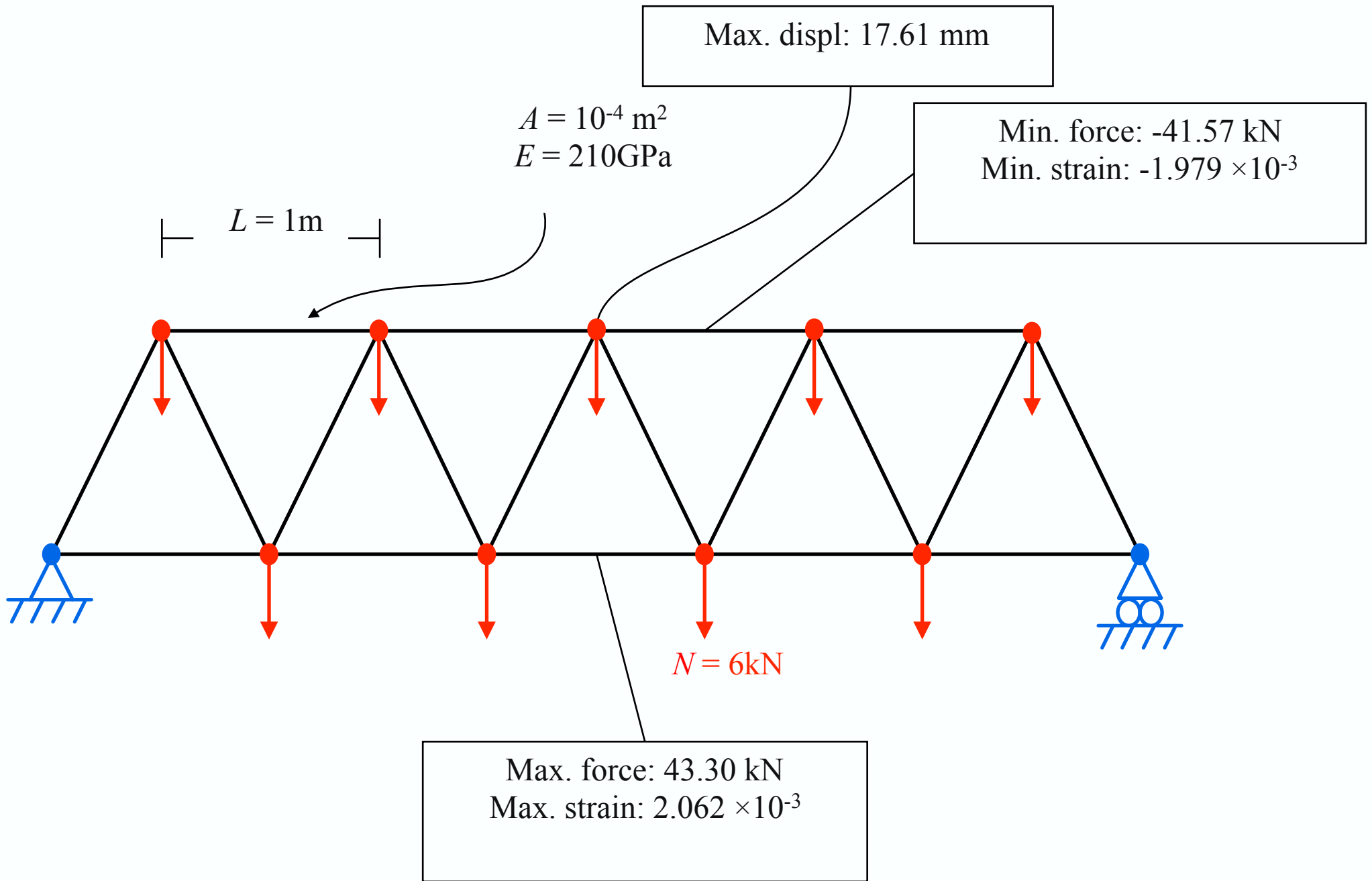
$$x(\xi) = \sum_{i=1}^{i=m} x_i N_i(\xi) , \quad x(\xi_i) = x_i ; \quad \text{coordinate at node } i$$

TRUSSES are structures made of bars connected to the rest of the structure by a **pin joint** at each end; a **joint** that can transmit forces but **cannot transmit moments**. As a result, moment equilibrium of the bar dictates that the bar is under **axial forces only**, since the bars have **no distributed loads**. **Loading** is only applied **at the nodes** of the structure.

Most engineering structures have a **skeleton** made of **long beams connected to each other** and on which a skin is sometimes added (roof structures, tubular design in cars, aircraft fuselage, satellites etc.) and sometimes not (bridges, grid transmission towers – even the Eiffel tower). A **truss approximation** of these structures (which assumes only axial forces in members and **forces on nodes** – but **no moments** – is a very **useful first approximation** on engineering to do **preliminary design** studies.

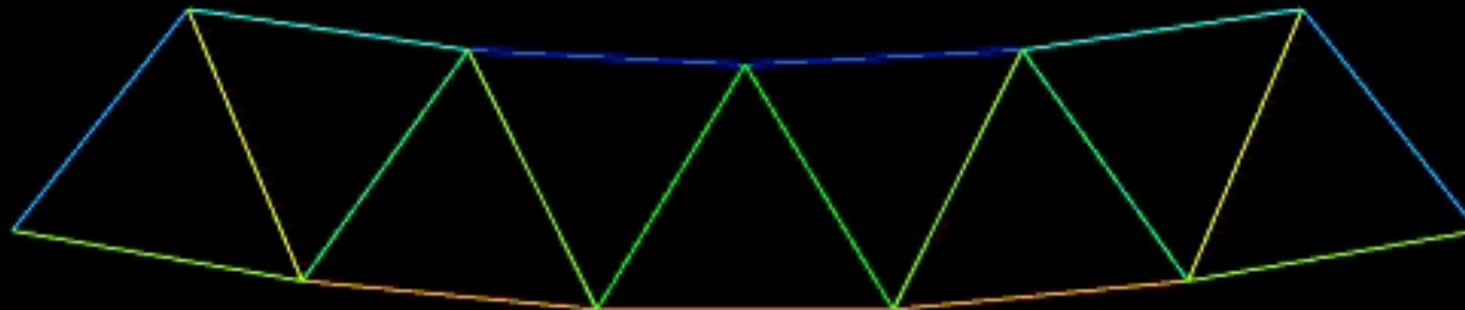
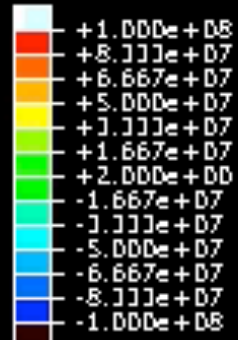
F.E.M. ideas apply here (linear elastic structure, small deformation and displacement linear kinematics) with each bar being an element.

TRUSS EXAMPLES IN 2D AND 3D



Step: Step-1 Frame: 100
Total Time: 1.000000

S, 511

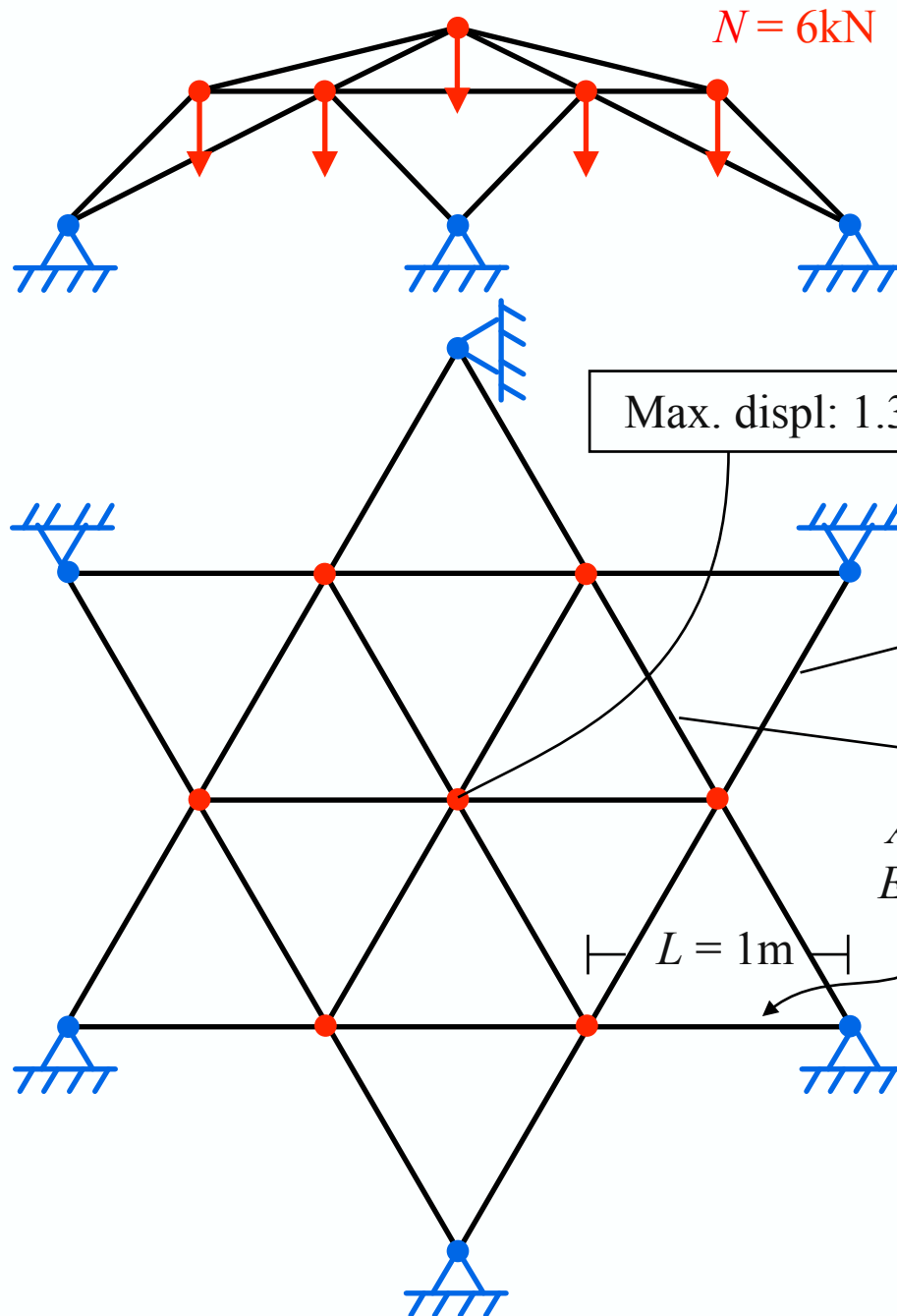


Y



ODB: Bridge.odb Abaqus/Standard 6.12-3 Mon Sep 15 16:16:58 GMT+02:00 2014

Step: Step-1
Increment: 100; Step Time = 1.000
Primary Var: S, 511
Deformed Var: U Deformation Scale Factor: +1.000e+02



Gravity loading of roof

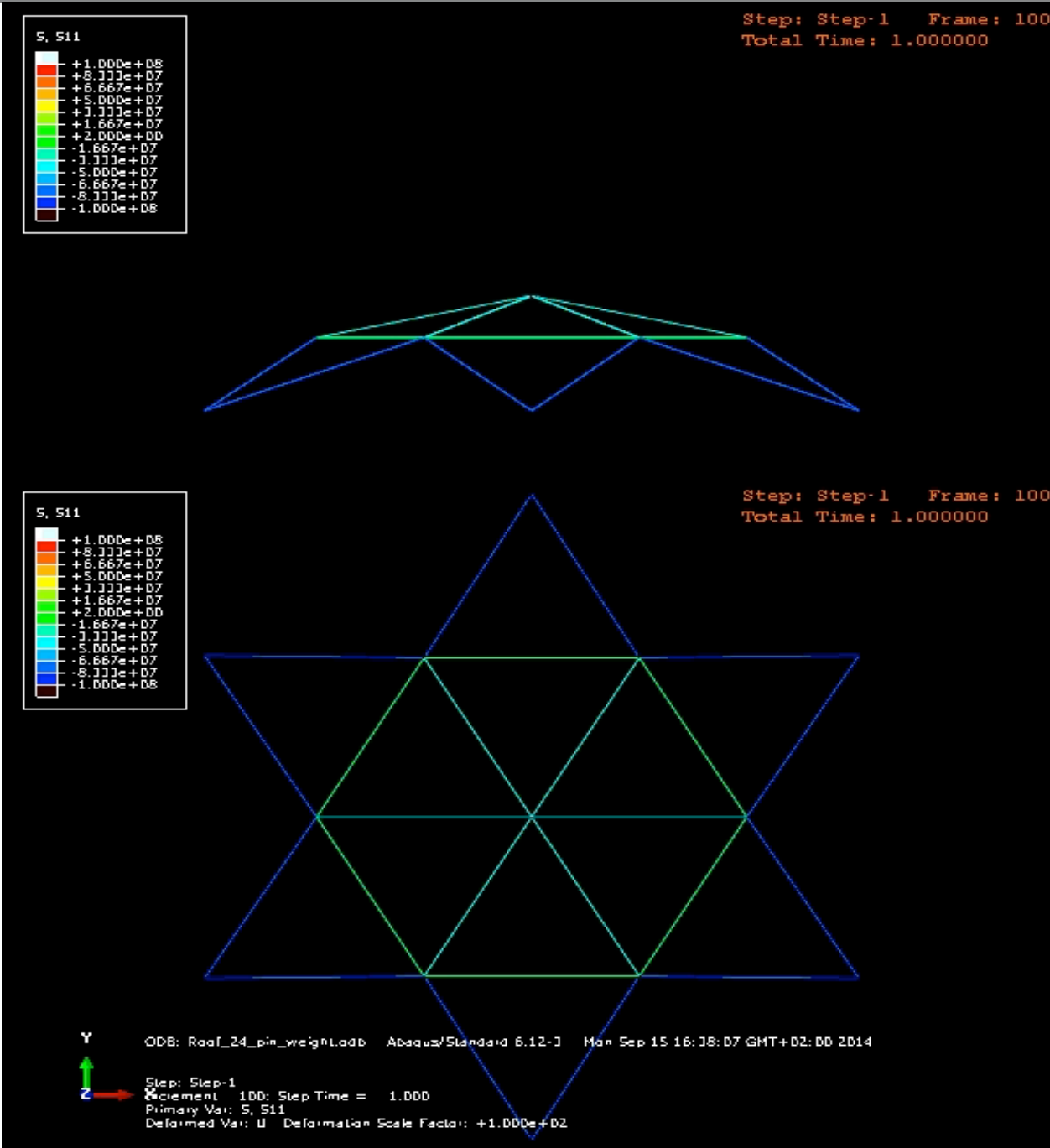
Max. displ: 1.338 mm

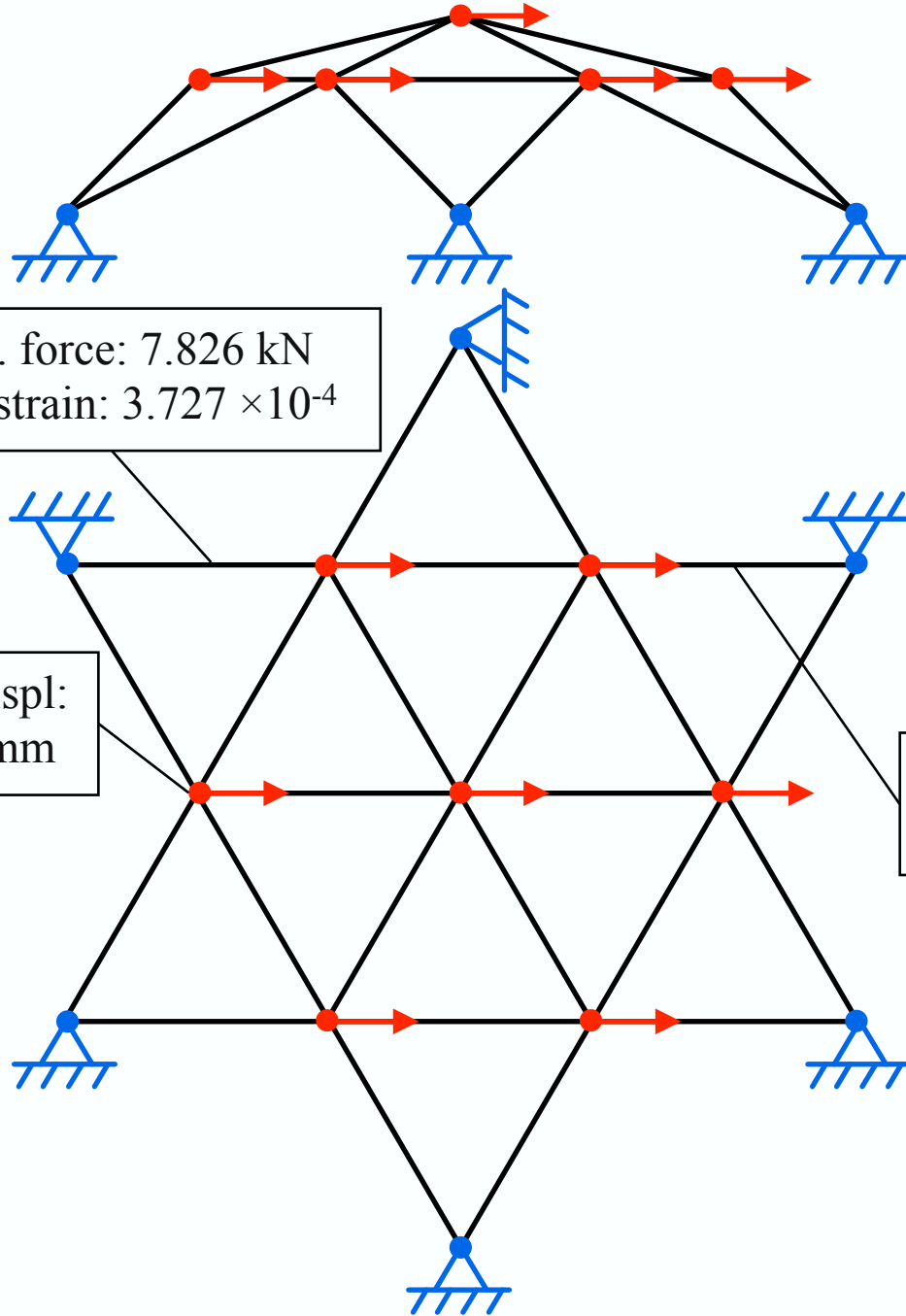
Min. force: -7.826 kN
Min. strain: -3.727×10^{-4}

Max. force: -3 kN
Max. strain: -1.429×10^{-4}

$A = 10^{-4} \text{ m}^2$
 $E = 210 \text{ GPa}$

$L = 1 \text{ m}$



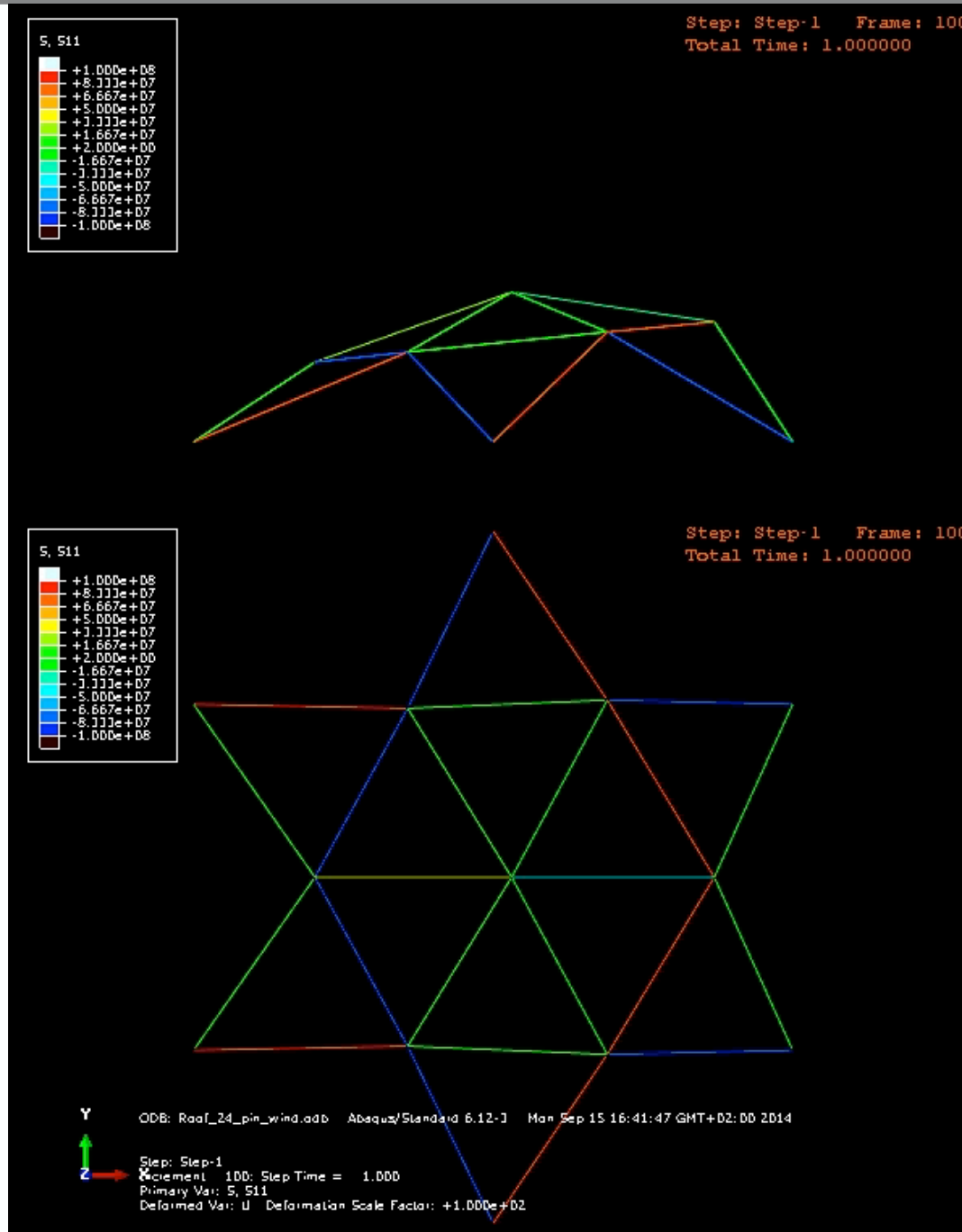


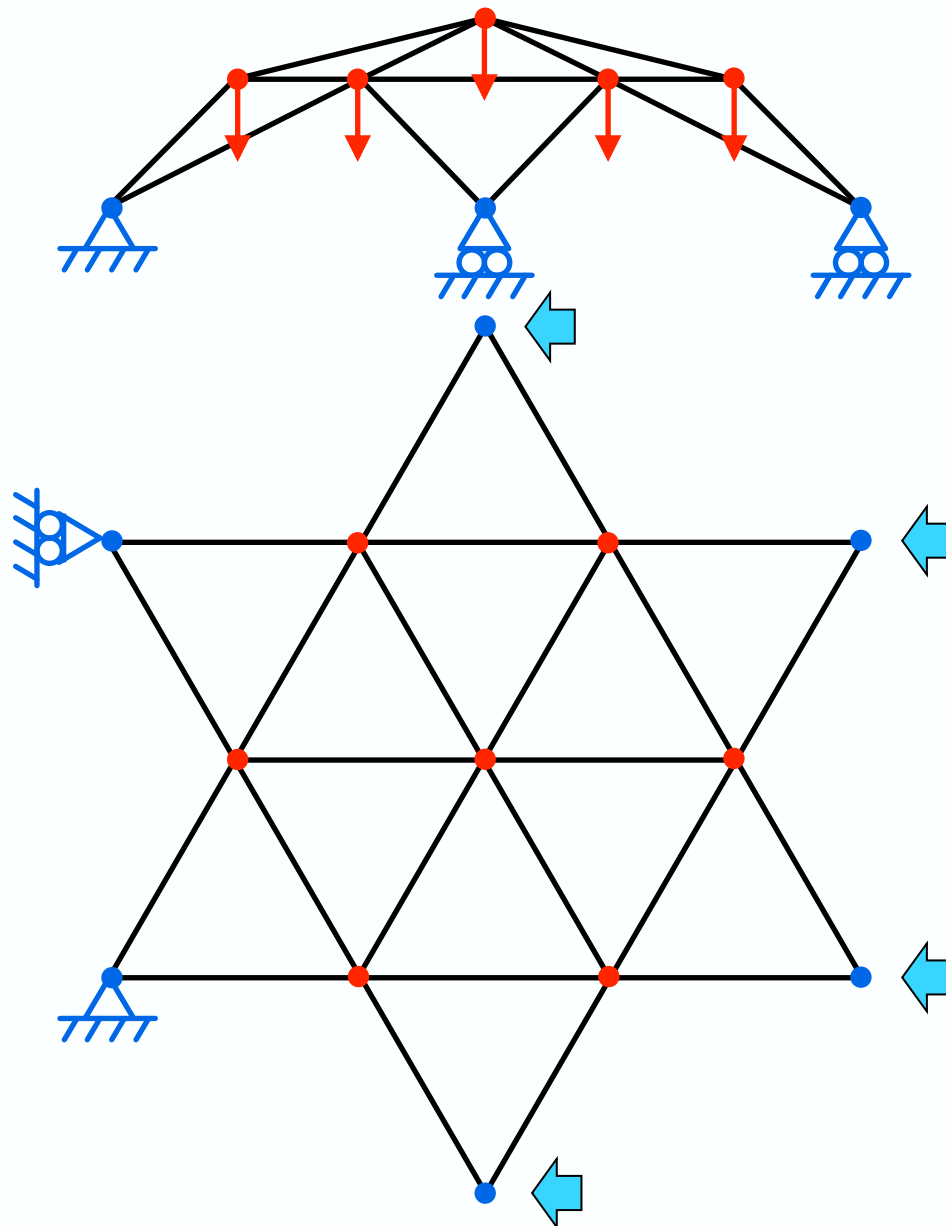
Wind loading of roof

Max. force: 7.826 kN
Max. strain: 3.727×10^{-4}

Max. displ:
1.474 mm

Min. force: -7.826 kN
Min. strain: -3.727×10^{-4}

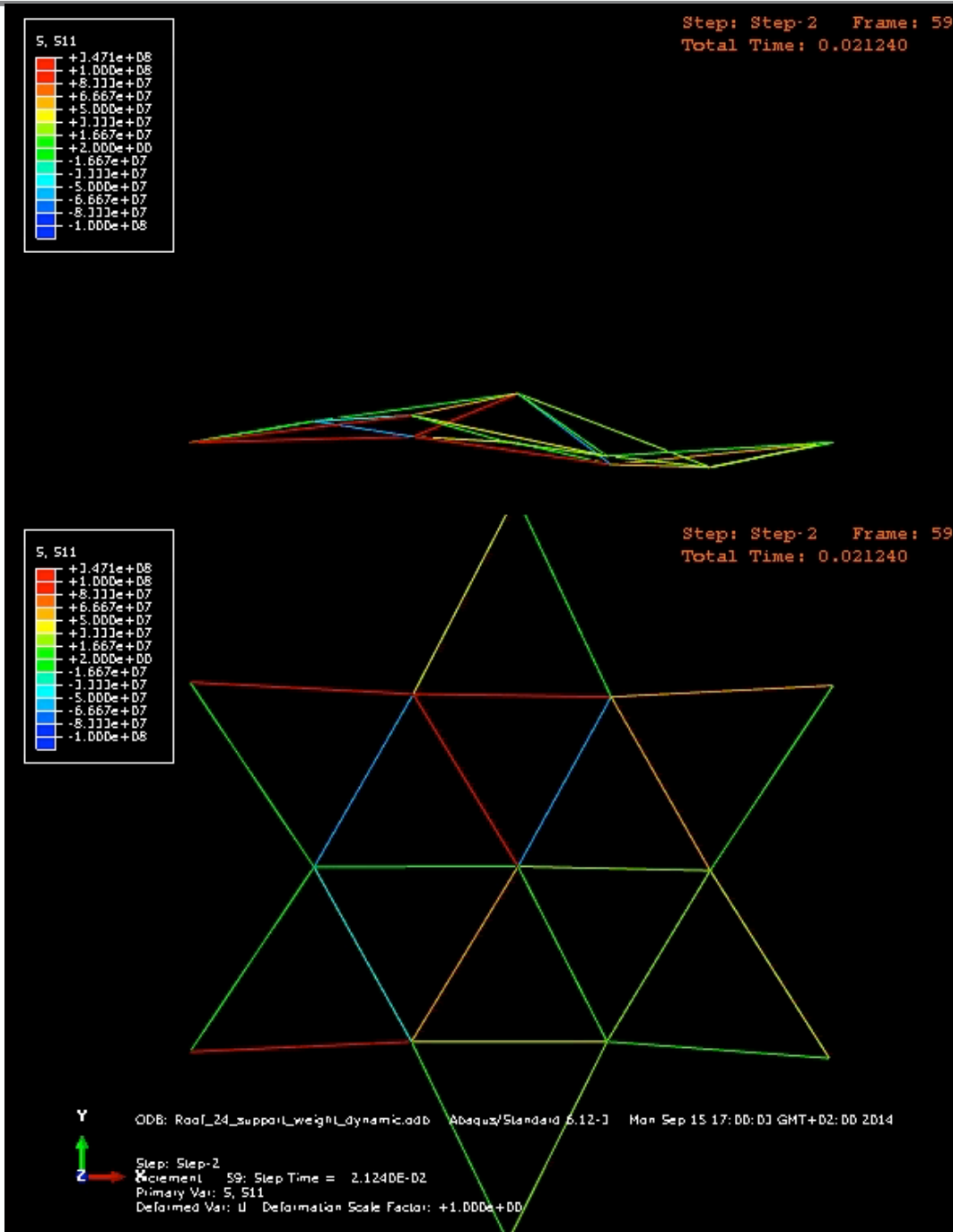




What happens when we change support conditions? (e.g. choose four ground nodes to be on rollers – no in-plane reactions)

Notice that the static problem is **ill conditioned** ($\det \mathbf{K} = 0$) and cannot be solved

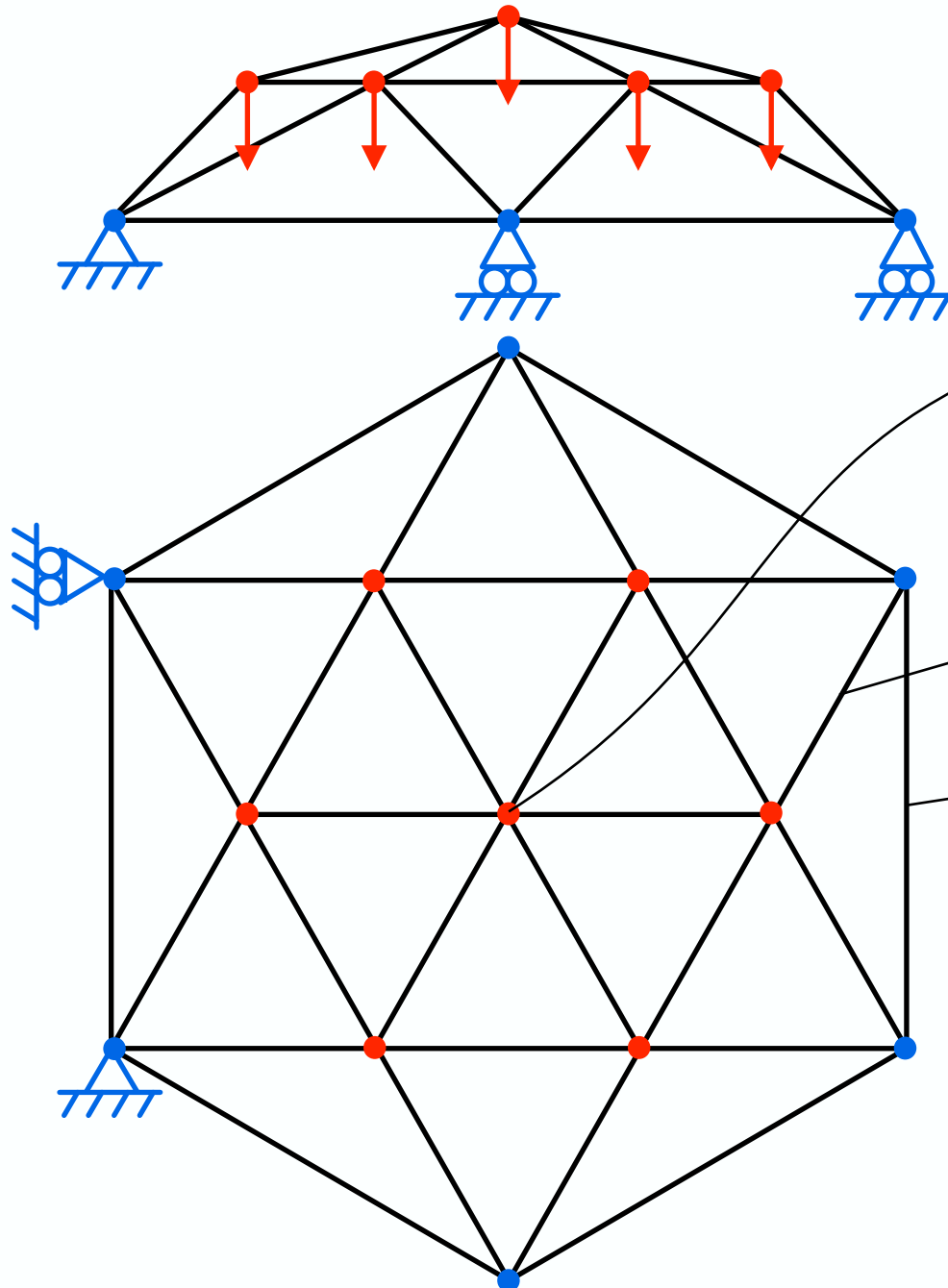
Structure is a mechanism! We can see what happens when we solve the dynamic problem...



A truss is a **mechanism** ($\det \mathbf{K} = 0$) when the number of available equilibrium equations exceeds the number of unknowns
(spatial dimension) \times (nodes) $>$ (bars) + (reactions)

A truss is a **isostatic** (also termed **statically determinate**) structure when the number of available equilibrium equations equals the number of unknowns; in this case you do not need the material properties of the bars, equilibrium equations suffice to solve the problem where bar forces that depend only on geometry!
(spatial dimension) \times (nodes) = (bars) + (reactions)

A truss is a **hyperstatic** (also termed **statically indeterminate**) structure when the number of available equilibrium equations is less than the number of unknowns; in this case bar forces depend on material properties
(spatial dimension) \times (nodes) $<$ (bars) + (reactions)



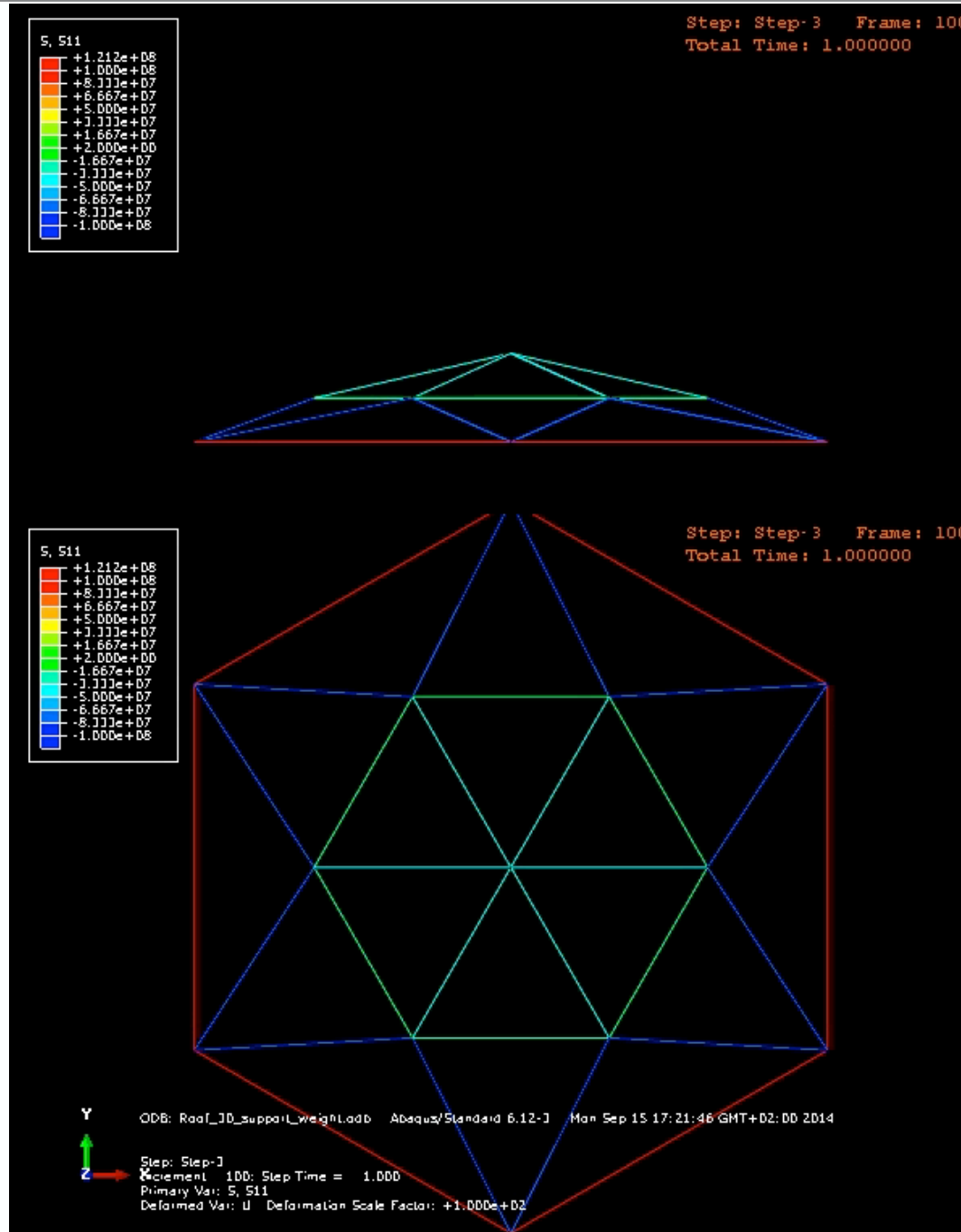
Gravity loading of roof
Modified structure that
allows roller supports

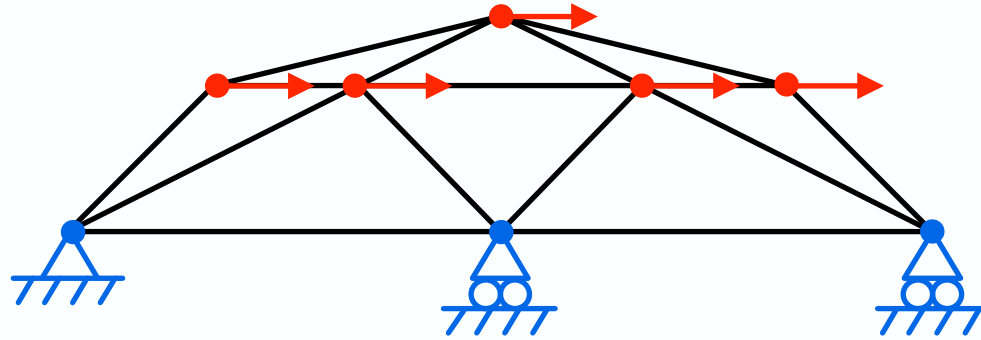
Max. displ: 3.228 mm

Min. force: -7.826 kN
 Min. strain: -3.727×10^{-4}

Max. force: 12.12 kN
 Max. strain: 5.774×10^{-4}

$(3) \times (13) = (30) + (9)$
Isostatic structure, forces
independent on bar material!

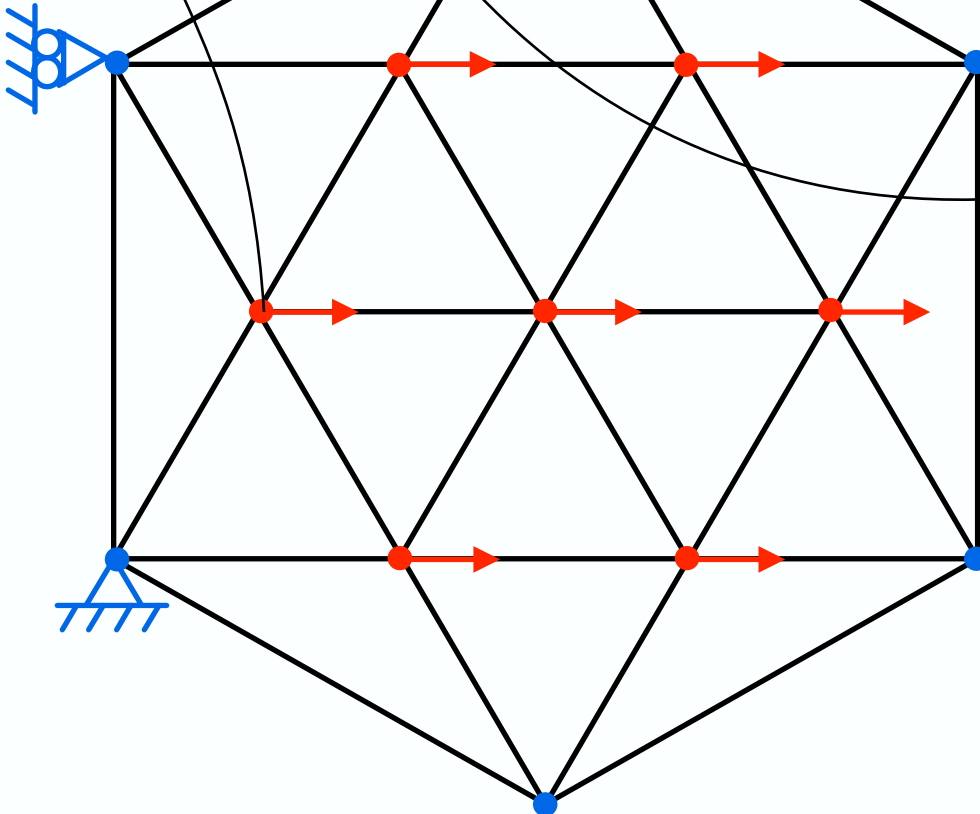




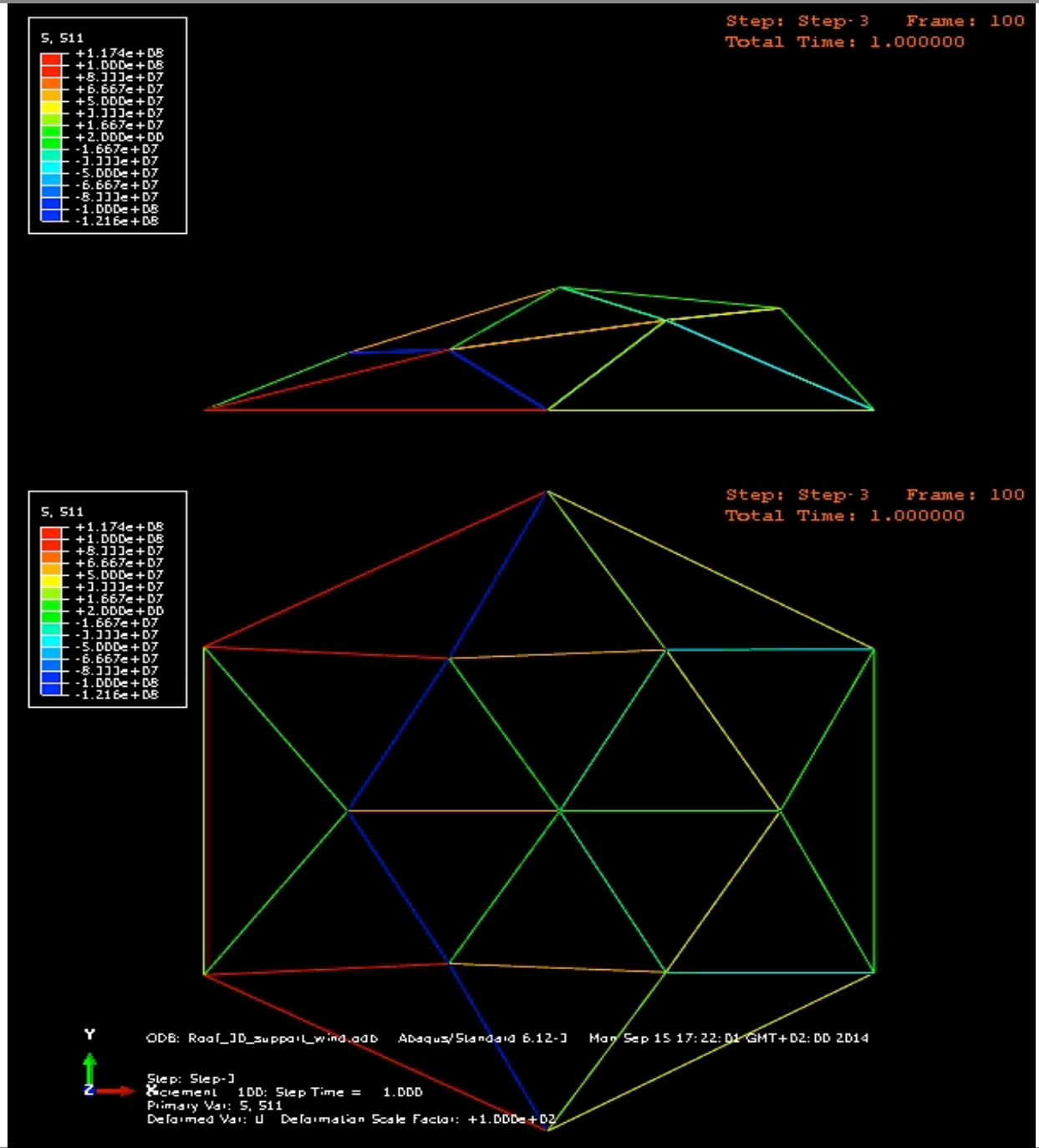
Wind loading of roof
Modified structure that
allows roller supports

Max. displ:
 2.625 mm

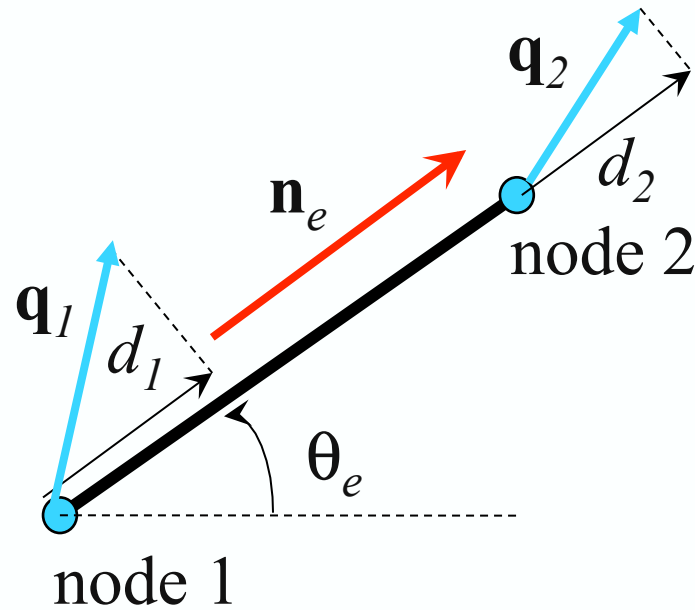
Max. force: 11.74 kN
 Max. strain: 5.590×10^{-4}



Min. force: -12.16 kN
 Min. strain: -5.790×10^{-4}



FORMULATION OF THE TRUSS PROBLEM



d_1, d_2 , axial displacements of bar

bar elongation: $d_2 - d_1$

d_i projection of displacement \mathbf{q}_i at node i

axial strain of element e : $\varepsilon_e = (d_2 - d_1)/l_e$

\mathbf{n}_e unit vector of element e

$$\mathbf{q}_e^T = [q_{1x}, q_{1y}, q_{2x}, q_{2y}], \quad \mathbf{n}_e = [n_{1x}, n_{1y}] = [\cos(\theta_e), \sin(\theta_e)] \quad \text{in 2D}$$

$$\mathbf{q}_e^T = [q_{1x}, q_{1y}, q_{1z}, q_{2x}, q_{2y}, q_{2z}], \quad \mathbf{n}_e = [n_{1x}, n_{1y}, n_{1z}] \quad \text{in 3D}$$

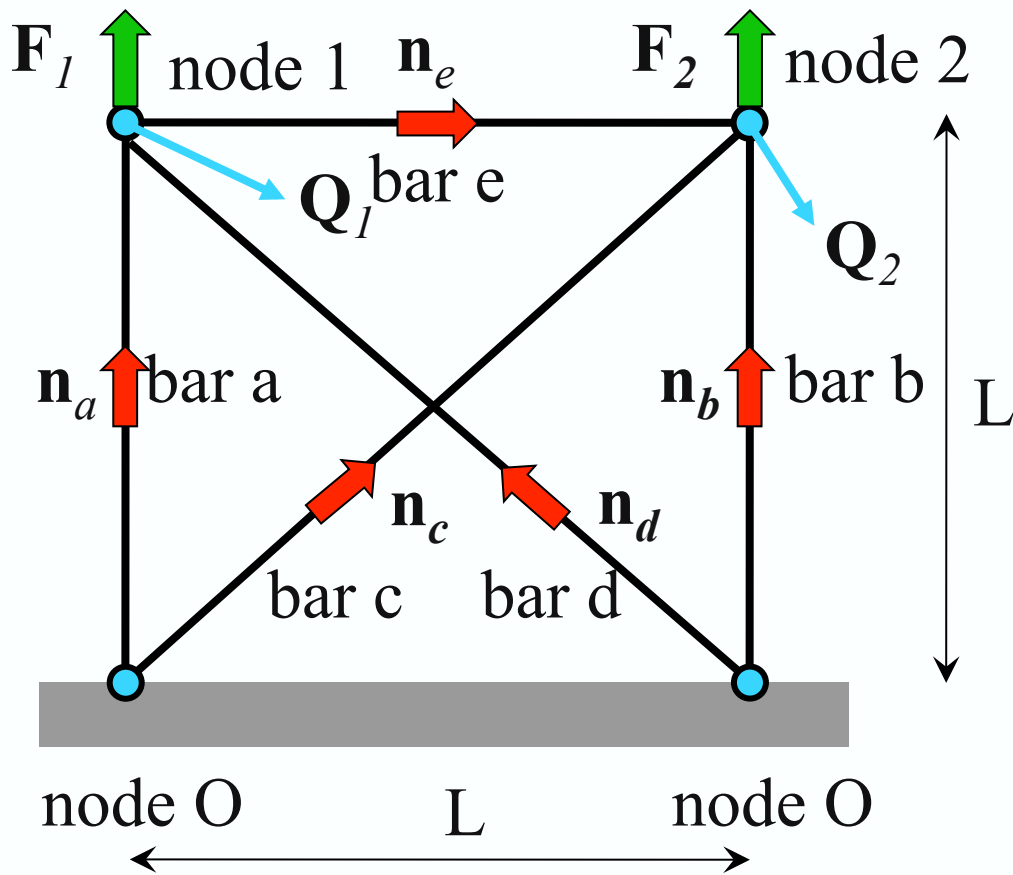
$$\mathcal{P} = \mathcal{P}_{int} + \mathcal{P}_{ext}$$

$$\mathcal{P}_{int} = \sum_e \mathcal{P}_{int}^e = \frac{1}{2} \mathbf{Q}^T \mathbf{K} \mathbf{Q} ; \quad \mathcal{P}_{int}^e = \int_{l_e} \left[\frac{1}{2} E A_e \epsilon_e^2 \right] dx = \frac{1}{2} \mathbf{q}_e^T \mathbf{k}_e \mathbf{q}_e$$

$$\epsilon_e = du/dx = (d_2 - d_1)/l_e ; \quad d_i = \mathbf{q}_i^T \mathbf{n}_e , \quad (i = 1, 2) , \quad \mathbf{q}_e^T = [\mathbf{q}_1^T , \mathbf{q}_2^T]$$

$$\mathbf{k}_e \equiv \begin{bmatrix} \mathbf{k}'_e & -\mathbf{k}'_e \\ -\mathbf{k}'_e & \mathbf{k}'_e \end{bmatrix} ; \quad \mathbf{k}'_e \equiv \mathbf{n}_e \mathbf{n}_e^T , \quad \mathbf{n}_e \text{ element orientation}$$

$$\mathcal{P}_{ext} = - \sum_{nodes} \mathbf{q}_n^T \mathbf{F}_n = \mathbf{Q}^T \mathbf{F} ; \quad \mathbf{q}_n , \mathbf{F}_n : \text{ displ., force at node } n$$



Uniform section bars (same EA)

$$\theta_a = \pi/2 \text{ for element a}$$

$$\theta_b = \pi/2 \text{ for element b}$$

$$\theta_c = \pi/4 \text{ for element c}$$

$$\theta_d = 3\pi/4 \text{ for element d}$$

$$\theta_e = 0 \text{ for element e}$$

$$\begin{bmatrix} \mathbf{F}_1 \\ \mathbf{F}_2 \end{bmatrix} = \begin{bmatrix} \mathbf{k}'_a + \mathbf{k}'_d + \mathbf{k}'_e & -\mathbf{k}'_e \\ -\mathbf{k}'_e & \mathbf{k}'_b + \mathbf{k}'_c + \mathbf{k}'_e \end{bmatrix} \begin{bmatrix} \mathbf{Q}_1 \\ \mathbf{Q}_2 \end{bmatrix}, \quad (\mathbf{F} = \mathbf{KQ})$$

$$\mathbf{k}'_e = \frac{EA}{l_e} \begin{bmatrix} \cos(\theta_e) \cos(\theta_e) & \cos(\theta_e) \sin(\theta_e) \\ \sin(\theta_e) \cos(\theta_e) & \sin(\theta_e) \sin(\theta_e) \end{bmatrix} \quad \text{for each element}$$

$$\mathbf{k}'_a = \mathbf{k}'_b = \frac{EA}{L} \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}, \quad \mathbf{k}'_e = \frac{EA}{L} \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix},$$

$$\mathbf{k}'_c = \frac{EA}{L\sqrt{2}} \begin{bmatrix} \frac{1}{2} & \frac{1}{2} \\ 1 & 1 \\ \frac{1}{2} & \frac{1}{2} \end{bmatrix}, \quad \mathbf{k}'_d = \frac{EA}{L\sqrt{2}} \begin{bmatrix} \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & \frac{1}{2} \end{bmatrix}$$

$$\begin{bmatrix} F_{1x} \\ F_{1y} \\ F_{2x} \\ F_{2y} \end{bmatrix} = \frac{EA}{L} \begin{bmatrix} 1 + \frac{1}{2\sqrt{2}} & -\frac{1}{2\sqrt{2}} & -1 & 0 \\ -\frac{1}{2\sqrt{2}} & 1 + \frac{1}{2\sqrt{2}} & 0 & 0 \\ -1 & 0 & 1 + \frac{1}{2\sqrt{2}} & \frac{1}{2\sqrt{2}} \\ 0 & 0 & \frac{1}{2\sqrt{2}} & 1 + \frac{1}{2\sqrt{2}} \end{bmatrix} \begin{bmatrix} Q_{1x} \\ Q_{1y} \\ Q_{2x} \\ Q_{2y} \end{bmatrix}$$

for loading : $\mathbf{F}^T = [0, 1, 0, 1]$

displacements are : $\mathbf{Q}^T = \frac{PL}{EA} \left[\frac{1}{3 + 4\sqrt{2}}, \frac{1 + 4\sqrt{2}}{3 + 4\sqrt{2}}, -\frac{1}{3 + 4\sqrt{2}}, \frac{1 + 4\sqrt{2}}{3 + 4\sqrt{2}} \right]$

Axial forces in each bar $N_e = EA(d_2 - d_1)/l_e$

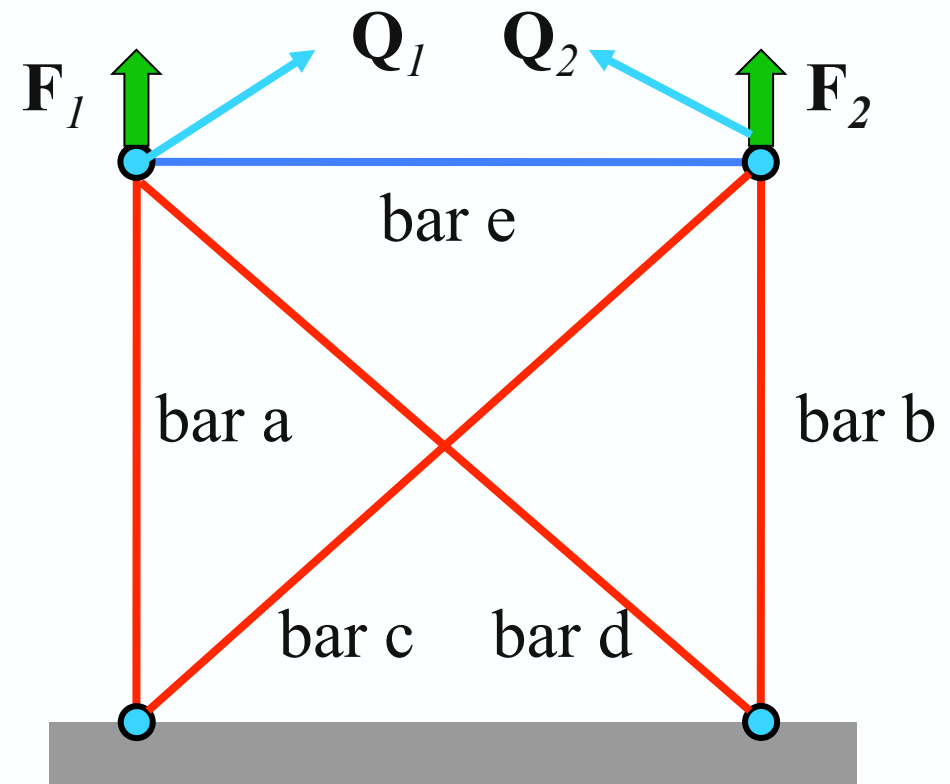
$$N_a = P (4\sqrt{2}+1)/(4\sqrt{2}+3) \text{ tensile}$$

$$N_b = P (4\sqrt{2}+1)/(4\sqrt{2}+3) \text{ tensile}$$

$$N_c = 2\sqrt{2} P / (4\sqrt{2}+3) \text{ tensile}$$

$$N_d = 2\sqrt{2} P / (4\sqrt{2}+3) \text{ tensile}$$

$$N_e = -2P / (4\sqrt{2}+3) \text{ compressive}$$



NOTE : Symmetric loading produces symmetric deformation & forces

NOTE : Taking into account thermal loading

$$\epsilon_e = du/dx - \alpha\Delta T = (d_2 - d_1)/l_e - \alpha\Delta T ; \quad d_i = \mathbf{q}_i^T \mathbf{n}_e$$

$$\mathcal{P}_{int}^e = \int_{l_e} \left[\frac{1}{2} E A_e \epsilon_e^2 \right] dx = \frac{1}{2} \mathbf{q}_e^T \mathbf{k}_e \mathbf{q}_e - \mathbf{q}_e^T \mathbf{f}_e$$

$$\mathbf{k}_e \equiv \begin{bmatrix} \mathbf{k}'_e & -\mathbf{k}'_e \\ -\mathbf{k}'_e & \mathbf{k}'_e \end{bmatrix} ; \quad \mathbf{k}'_e \equiv \mathbf{n}_e \mathbf{n}_e^T , \quad \mathbf{n}_e \text{ element orientation}$$

$$\mathbf{f}_e^T \equiv E A_e \alpha \Delta T \left[-\mathbf{n}_e^T , \mathbf{n}_e^T \right]$$

α : thermal expansion coefficient, ΔT : temperature change

For students in **PA in Mechanics (MEC 592, MEC 595)**, I will be at your disposal on **October 05** at **11:am** in **Amphi MONGE** to talk to you about internships in **Mechanics**.

I will also inform you about a dual MS degree program with Caltech (Departement de Mecanique at X and Aerospace Engineering at Caltech) which concerns students interested in pursuing a Doctorate degree in either Fluid Mechanics or Solid Mechanics back here (LadHyx or LMS) and talk about similar possibilities exist with University of Minnesota for Aerospace and Civil Engineering (LMS)